



The First Negative-order Matrix AKNS Flow and its Darboux Transformation

Wen-Xiu Ma^{1,2,3,4}

Received: 26 May 2025 / Accepted: 4 August 2025

© The Author(s), under exclusive licence to Springer Science+Business Media, LLC, part of Springer Nature 2025

Abstract

This paper presents the first negative-order matrix AKNS flow, derived by associating a Lax pair containing a first-order pole in the spectral parameter with the matrix AKNS spectral problem. The corresponding Darboux transformation is constructed within the AKNS framework. Starting from a seed solution, a class of exact and explicit solutions to the nonlinear negative-order model is generated via a single application of the derived Darboux transformation.

Keywords Lax pair · Matrix AKNS flows · Darboux transformation

Mathematics Subject Classification (2010) 37K10 · 35Q51 · 37K40 · 02.30.Ik · 05.45.Yv

1 Introduction

The Darboux transformation is a powerful algebraic method for generating new solutions from known ones in integrable systems, especially in soliton theory [1]. Applied to the Lax pair of an integrable system, it preserves integrability and enables the construction of multi-soliton solutions, rational solutions, rogue waves, and other complex structures. The method has been extended to matrix, higher-dimensional, noncommutative, and supersymmetric settings (see, e.g., [2–4]). Consider the zero-curvature formulation for integrable systems, in which a system of PDEs

✉ Wen-Xiu Ma
mawx@cas.usf.edu

¹ Present address: Department of Mathematics, Zhejiang Normal University, Jinhua 321004, China

² Research Center of Astrophysics and Cosmology, Khazar University, Baku, Azerbaijan

³ Department of Mathematics and Statistics, University of South Florida, 4202 E Fowler Avenue, Tampa, FL, USA

⁴ Material Science Innovation and Modelling, North-West University, Mafikeng Campus, Mafikeng, South Africa

$$u_t = K(u) = K(x, t, u, u_x, \dots), \quad (1)$$

is represented by the zero-curvature condition:

$$U_t - V_x + [U, V] = 0, \quad (2)$$

where U and V are matrix-valued functions (the Lax pair) belonging to a matrix loop algebra [5]. This zero-curvature equation arises as the compatibility condition of the matrix spectral problems

$$\phi_x = U\phi = U(u, \lambda)\phi, \quad \phi_t = V\phi = V(u, \lambda)\phi, \quad (3)$$

where λ is the spectral parameter and ϕ is the vector eigenfunction.

A Darboux transformation consists of a gauge transformation $\phi' = D\phi$ and a new potential $u' = u'(u)$, with $D = D(u, \lambda)$ a matrix function of λ , such that the transformed function ϕ' solves a pair of matrix spectral problems of the same structure:

$$\phi'_x = U'\phi' = U(u', \lambda)\phi', \quad \phi'_t = V'\phi' = V(u', \lambda)\phi', \quad (4)$$

where U' and V' retain the structure of the original Lax pair. The Darboux matrix D must satisfy the two compatibility conditions:

$$U'D = DU + D_x, \quad V'D = DV + D_t. \quad (5)$$

One of important tasks in the field of integrable systems is to construct such Darboux transformations explicitly. Suppose that U and V are $N \times N$ matrices. A first-order transformation can often be chosen in the form:

$$D(\lambda) = \lambda I_N - S, \quad (6)$$

where I_N is the identity matrix of order N and S is an $N \times N$ matrix independent of λ . Given N eigenvalues $\lambda_1, \lambda_2, \dots, \lambda_N$, let $\phi^{[j]}$ denote the corresponding eigenfunctions that satisfy:

$$\phi_x^{[j]} = U(u, \lambda_j)\phi^{[j]}, \quad \phi_t^{[j]} = V(u, \lambda_j)\phi^{[j]}, \quad 1 \leq j \leq N, \quad (7)$$

for some fixed solution u of (1). Then, a standard choice for the matrix S is:

$$S = HAH^{-1}, \quad H = (\phi^{[1]}, \dots, \phi^{[N]}), \quad (8)$$

where $A = \text{diag}(\lambda_1, \dots, \lambda_N)$, as described in [2, 6]. Verifying the compatibility conditions (5) ensures the validity of the transformation.

In this letter, we begin with the matrix AKNS spectral problem and propose the first negative-order matrix AKNS flow. We construct its associated Darboux transformation within the AKNS framework and apply it once to generate a class of explicit solutions. The final section offers concluding remarks and discusses directions for future research.

2 The Lax Pair and the Matrix Model

Let α_1 and α_2 be two distinct constants, and let m and n be natural numbers. We consider the AKNS spectral matrix [7] of the form:

$$U = i\lambda\Lambda + Q, \quad \Lambda = \begin{bmatrix} \alpha_1 I_m & 0 \\ 0 & \alpha_2 I_n \end{bmatrix}, \quad Q = \begin{bmatrix} 0 & q \\ r & 0 \end{bmatrix}, \quad (9)$$

where q and r are the potential matrices of sizes $m \times n$ and $n \times m$, respectively. We define the Lax operator by

$$V = -\frac{1}{\alpha\lambda}W, \quad \alpha = \alpha_1 - \alpha_2, \quad (10)$$

where

$$W = W_1 + W_2, \quad W_1 = -i \begin{bmatrix} v & 0 \\ 0 & w \end{bmatrix}, \quad W_2 = i \begin{bmatrix} 0 & -q_t \\ r_t & 0 \end{bmatrix}, \quad (11)$$

with v and w being two square potential matrices of orders m and n , respectively.

The zero-curvature equation

$$U_t - V_x + [U, V] = 0 \quad (12)$$

is equivalent to

$$-\alpha\lambda U_t - W_x + [U, W] = 0. \quad (13)$$

By equating coefficients of powers of λ , we obtain the system:

$$\begin{cases} -\alpha Q_t + i[\Lambda, W] = 0, \\ -W_x + [Q, W] = 0. \end{cases} \quad (14)$$

The first equation is satisfied identically due to the choices:

$$[\Lambda, W_1] = 0, \quad [\Lambda, W_2] = -i\alpha Q_t \quad (15)$$

For the second equation, we compute

$$[Q, W_1] = -i \begin{bmatrix} 0 & qw - vq \\ rv - wr & 0 \end{bmatrix}, \quad [Q, W_2] = i \begin{bmatrix} (qr)_t & 0 \\ 0 & -(rq)_t \end{bmatrix}. \quad (16)$$

Thus, the second equation yields:

$$\begin{cases} q_{tx} - qw + vq = 0, \\ r_{tx} + rv - wr = 0, \\ v_x + (qr)_t = 0, \\ w_x - (rq)_t = 0. \end{cases} \quad (17)$$

Since V contains a simple pole of λ , this system represents the first negative-order matrix AKNS flow. It should be noted that this represents a special case of the system introduced in [6], corresponding to the parameter choices

$$n = 2, \quad p_1 = x, \quad x_2 = t, \quad A_1 = -U, \quad a_2 = -\alpha\lambda, \quad A_2 = W. \quad (18)$$

All functions involved in the Lax pair U and V , including q and r , depends solely on x and t . Consequently, there is no dependence on x_1 and p_2 .

If we take

$$v = -\frac{a}{2} - \partial_x^{-1}(qr)_t, \quad w = \frac{a}{2} + \partial_x^{-1}(rq)_t, \quad (19)$$

where a is an arbitrary constant, then the system simplifies to:

$$\begin{cases} q_{tx} - q\partial_x^{-1}(rq)_t - \partial_x^{-1}(qr)_t q = aq, \\ r_{tx} - r\partial_x^{-1}(qr)_t - \partial_x^{-1}(rq)_t r = ar. \end{cases} \quad (20)$$

When $a = 1$, this system corresponds to the first negative-order AKNS system studied in [8]. Furthermore, in the scalar case $m = n = 1$, it reduces to the first negative-order scalar AKNS system presented in [9, 10]. There are other negative-order AKNS flows (see, e.g., [11–13]). It is worth noting that the resulting equations do not depend explicitly on the constants α_1 and α_2 .

Let

$$m = n = 1, \quad w = -v. \quad (21)$$

Then, the system becomes

$$\begin{cases} q_{tx} + 2vq = 0, \\ r_{tx} + 2vr = 0, \\ v_x + (qr)_t = 0. \end{cases} \quad (22)$$

With $r = q$, this reduces to

$$\begin{cases} q_{tx} + 2vq = 0, \\ v_x + (q^2)_t = 0. \end{cases} \quad (23)$$

Let

$$m = n = 1, \quad r = q(-x, t), \quad w = v(-x, t), \quad (24)$$

then the system becomes a nonlocal model:

$$\begin{cases} q_{tx} + (v - v(-x, t))q = 0, \\ v_x + (qq(-x, t))_t = 0. \end{cases} \quad (25)$$

3 Darboux Transformation

3.1 Compatibility Conditions

We assume that the Darboux matrix takes the form:

$$D(\lambda) = \lambda I_{m+n} - S. \quad (26)$$

To satisfy the Darboux transformation conditions, this matrix must fulfill the spatial requirement

$$U'D = DU + D_x, \quad (27)$$

where U is given by (9), and the transformed operator is

$$U' = i\lambda\Lambda + Q'. \quad (28)$$

Substituting into the requirement yields:

$$(i\lambda\Lambda + Q')(\lambda I_{m+n} - S) = (\lambda I_{m+n} - S)(i\lambda\Lambda + Q) - S_x,$$

which implies:

$$Q' = Q + i[\Lambda, S], \quad (29)$$

and

$$S_x = Q'S - SQ = [Q + i\Lambda S, S]. \quad (30)$$

Next, we consider the temporal compatibility condition:

$$V'D = DV + D_t, \quad (31)$$

where

$$V' = -\frac{1}{\alpha\lambda}W'. \quad (32)$$

This gives:

$$W'(\lambda I_{m+n} - S) = (\lambda I_{m+n} - S)W + \alpha \lambda S_t,$$

which implies:

$$W' = W + \alpha S_t, \quad W'S - SW = 0.$$

Thus, we have

$$W' = SW S^{-1}, \quad (33)$$

and

$$S_t = \frac{1}{\alpha}(SW S^{-1} - W). \quad (34)$$

3.2 Construction of the Darboux Matrix

Following the general framework (see, e.g., [2, 6]), we formulate the matrix S as:

$$S = HAH^{-1}, \quad (35)$$

where

$$H = (\phi^{[1]}, \dots, \phi^{[m+n]}), \quad A = \text{diag}(\lambda_1, \dots, \lambda_{m+n}), \quad (36)$$

and each column vector $\phi^{[j]}$ satisfies the matrix spectral problems:

$$\phi_x^{[j]} = U(u, \lambda_j)\phi^{[j]}, \quad \phi_t^{[j]} = V(u, \lambda_j)\phi^{[j]}, \quad 1 \leq j \leq m+n. \quad (37)$$

It follows that

$$H_x = i\Lambda H A + QH, \quad H_t = WHB, \quad (38)$$

where

$$B = \text{diag}\left(-\frac{1}{\alpha\lambda_1}, \dots, -\frac{1}{\alpha\lambda_{m+n}}\right). \quad (39)$$

We can then compute:

$$\begin{aligned}
S_x &= H_x A H^{-1} - H A (H^{-1} H_x H^{-1}) \\
&= i \Lambda H A^2 H^{-1} + Q H A H^{-1} - H A H^{-1} (i \Lambda H A + Q H) H^{-1} \\
&= i \Lambda H A^2 H^{-1} + Q H A H^{-1} - i H A H^{-1} \Lambda H A H^{-1} - H A H^{-1} Q \\
&= Q S - S Q + i \Lambda S^2 - i S \Lambda S,
\end{aligned}$$

which confirms that the spatial condition (30) is satisfied.

Similarly, from the second equality in (38), we obtain:

$$S_t = H_t A H^{-1} - H A (H^{-1} H_t H^{-1}) = W H B A H^{-1} - H A (H^{-1} W H B H^{-1}).$$

Using the identity $\alpha B A = -I_{m+n}$, it simplifies to:

$$\begin{aligned}
\alpha S_t &= W H (\alpha B A) H^{-1} - H A H^{-1} W H (\alpha B) H^{-1} \\
&= -W H H^{-1} + H A H^{-1} W H A^{-1} H^{-1} \\
&= -W + S W S^{-1} = -W + W',
\end{aligned}$$

confirming the temporal condition (34).

3.3 The Darboux Transformation

We thus obtain the Darboux transformation:

$$\phi' = (\lambda I_{m+n} - S)\phi, \quad U' = i\lambda\Lambda + Q', \quad V' = -\frac{1}{\alpha\lambda}W', \quad (40)$$

with

$$Q' = Q + i[\Lambda, S], \quad W' = S W S^{-1}, \quad (41)$$

where S is given by (35). Based on the structure of U and W in (9) and (11), this transformation gives the following expressions for the transformed quantities:

$$\begin{cases} q' = q + i[\Lambda, S]_{12}, & r' = r + i[\Lambda, S]_{21}, \\ v' = i(S W S^{-1})_{11}, & w' = i(S W S^{-1})_{22}, \end{cases} \quad (42)$$

where M_{jk} denotes the (j, k) -th block of the matrix M .

3.4 Explicit Solutions

As an application, consider the seed solution

$$q = r = 0, \quad v = v_0, \quad w = w_0, \quad (43)$$

where v_0 and w_0 are constant matrices. Then, the eigenfunctions take the form:

$$\phi^{[j]} = (\phi_1^{[j]T}, \phi_2^{[j]T})^T, \quad (44)$$

with

$$\begin{cases} \phi_1^{[j]} = \exp(i\alpha_1 \lambda_j I_m x + \frac{i}{\alpha \lambda_j} v_0 t) \mu_1^{[j]}, \\ \phi_2^{[j]} = \exp(i\alpha_2 \lambda_j I_n x + \frac{i}{\alpha \lambda_j} w_0 t) \mu_2^{[j]}, \end{cases} \quad 1 \leq j \leq m+n, \quad (45)$$

where $\mu_1^{[j]}$ and $\mu_2^{[j]}$ are arbitrary constant vectors of sizes m and n , respectively. Then, the following explicit solutions to the system (17) are obtained:

$$q' = i[\Lambda, S]_{12}, \quad r' = i[\Lambda, S]_{21}, \quad v' = i\alpha(SWS^{-1})_{11}, \quad w' = i(SWS^{-1})_{22}, \quad (46)$$

where $S = HAH^{-1}$ with $H = (\phi^{[1]}, \dots, \phi^{[m+n]})$ and $A = \text{diag}(\lambda_1, \dots, \lambda_{m+n})$.

4 Concluding Remarks

We have constructed the first negative-order matrix AKNS flow and its associated Darboux transformation, along with a class of explicit solutions generated from a single application of the transformation. Iterating this process N times yields a higher-order Darboux transformation. Other developments on Darboux transformations for negative-order flows can be found in the literature (see, e.g., [14–16]).

An intriguing question is whether negative-order flows form commuting hierarchies, analogous to their well-known positive-order counterparts. A broader generalization involves Lax operators with higher-order poles in the spectral parameter, and a major open problem is to characterize Darboux transformations in this extended, higher negative-order framework.

Constrained Lax pairs are of particular interest, especially when reductions are imposed on the general spectral matrix U (see, e.g., [17, 18]). Such reductions often lead to nonlocal integrable models, which themselves form commuting hierarchies [19]. A central challenge is to develop systematic methods for constructing Darboux transformations in these constrained matrix settings.

Acknowledgements The work was supported in part by the Ministry of Science and Technology of China (G2021016032L and G2023016011L) and the National Natural Science Foundation of China (12271488 and 11975145). The author would also like to thank Dr. Zeeshan Amjad, Dr. Asifa Ashraf, Dr. Aamir Farooq and Dr. Saira Hameed for their valuable discussions.

Author Contributions WX wrote the main manuscript text, made the computations, and reviewed the manuscript.

Data Availability No datasets were generated or analysed during the current study.

Declarations

Ethical Approval Not applicable.

Competing interests The author declares no competing interests.

References

1. Matveev, V.B., Salle, M.A.: *Darboux Transformation and Solitons*. Springer, Berlin (1991)
2. Gu, C.H., Hu, H.S., Zhou, Z.X.: *Darboux Transformations in Integrable Systems*. Springer, Dordrecht (2005)
3. Doktorov, E.V., Leble, S.B.: *A Dressing Method in Mathematical Physics*. Springer, Dordrecht (2007)
4. Ma, W.X., Zhang, Y.J.: Darboux transformations of integrable couplings and applications. *Rev. Math. Phys.* **30**(2), 1850003 (2018)
5. Ma, W.X., Chen, M.: Hamiltonian and quasi-Hamiltonian structures associated with semi-direct sums of Lie algebras. *J. Phys. A: Math. Gen.* **39**(34), 10787–10801 (2006)
6. Ma, W.X.: Darboux transformations for a Lax integrable system in $2n$ -dimensions. *Lett. Math. Phys.* **39**(1), 33–49 (1997)
7. Ablowitz, M.J., Kaup, D.J., Newell, A.C., Segur, H.: The inverse scattering transform-Fourier analysis for nonlinear problems. *Stud. Appl. Math.* **53**(4), 249–315 (1974)
8. Chen, K., Liu, S.M., Zhang, D.J.: Covariant hodograph transformations between nonlocal short pulse models and the AKNS(-1) system. *Appl. Math. Lett.* **88**, 230–236 (2019)
9. Zhang, D.J., Ji, J., Zhao, S.L.: Soliton scattering with amplitude changes of a negative order AKNS equation. *Phys. D* **238**(23–24), 2361–2367 (2009)
10. Amjad, Z., Khan, D.: Binary Darboux transformation for a negative-order AKNS equation. *Theor. Math. Phys.* **206**(2), 128–141 (2021)
11. Ji, J., Zhang, J.B., Zhang, D.J.: Soliton solutions for a negative order AKNS equation hierarchy. *Commun. Theor. Phys.* **52**(3), 395–397 (2009)
12. Yao, Y.Q., Zeng, Y.B.: Coupled short pulse hierarchy and its Hamiltonian structure. *J. Phys. Soc. Jpn.* **80**(6), 064004 (2011)
13. Yu, G.F., Lao, D.: Complex and coupled complex negative order AKNS equation. *Commun. Nonlinear Sci. Numer. Simul.* **30**(1–3), 196–206 (2016)
14. Chen, C., Zhou, Z.X.: Darboux transformation and exact solutions of the Myrzakulov-I equation. *Chin. Phys. Lett.* **26**(8), 080504 (2009)
15. Riaz, H., Wajahat, A., ul Hassan, M.: Noncommutative negative order AKNS equation and its soliton solutions. *Mod. Phys. Lett. A* **33**(35), 1850209 (2018)
16. Wajahat, H., Riaz, A.: Darboux transformation for a negative order AKNS equation. *Commun. Theor. Phys.* **71**(8), 912–920 (2019)
17. Ma, W.X.: Integrable matrix nonlinear Schrödinger equations with reduced Lax pairs of AKNS type. *Appl. Math. Lett.* **168**, 109574 (2025)
18. Ma, W.X.: Matrix mKdV integrable hierarchies via two identical group reductions. *Mathematics* **13**(9), 1438 (2025)
19. Ma, W.X.: Nonlocal integrable equations in soliton theory. in: *Nonlinear and Modern Mathematical Physics*, eds. Manukure, S., Ma, W.X., Springer, Cham, Switzerland, pp. 251–266. (2024)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Springer Nature or its licensor (e.g. a society or other partner) holds exclusive rights to this article under a publishing agreement with the author(s) or other rightsholder(s); author self-archiving of the accepted manuscript version of this article is solely governed by the terms of such publishing agreement and applicable law.