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# Frobenius integrable decompositions for PDEs

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# Outline

**1. Introduction**

**2. Specific PDEs possessing**

**Frobenius integrable decompositions (FIDs)**

**3. Conclusions**



# 1. Introduction

- Many PDEs exhibiting soliton phenomena appeared in many science subjects —— *fluid physics, solid physics, elementary particle physics, biological physics, superconductor physics, ...*
- It is a quite fascinating research topic that how to solve such PDEs to obtain interesting solutions including *solitons*, which attracts much attention of mathematicians, physicists and dynamicists.



- In the past several decades, there have been many efficient methods for constructing exact solutions to nonlinear PDEs.
- Though the solving methods are diverse, appropriate **reductions**
  - *similarity reductions,*
  - *symmetry constraints,*
  - *travelling wave reductions*
  -
- To reduce given PDEs to **simpler** PDEs and/or integrable ODEs.



- W.X. Ma, H.Y. Wu, J.S. He, Partial differential equations possessing Frobenius integrable decompositions, *Physics Letters A*, 364 (2007) 29-32.
- They presented **FIDs** for two classes of nonlinear evolution equations (NEEs) with logarithmic derivative Bäcklund transformations in soliton theory.

$$u = (\ln \phi)_x = \frac{\psi}{\phi}, \quad u = (\ln \phi)_{xx} = \frac{\lambda \phi^2 - \psi^2}{\phi^2}.$$

- The discussed NEEs are transformed into systems of Frobenius integrable ODEs with cubic nonlinearity.



- F.C. You, T.C. Xia, J. Zhang, Frobenius integrable decompositions for two classes of nonlinear evolution equations with variable coefficients, *Modern Physics Letters B*, 23 (12) (2009) 1519-1524
- They obtained two classes of PDEs with variable coefficients possessing **FIDs**, including
  - the KdV equation
  - the potential KdV equation
  - the Boussinesq equation
  - the generalized BBM equation, ...



## What is FIDs ?

$$P(u, u_t, u_x, u_{xt}, \dots) = 0$$

$$u = v(\Phi) = v(\Phi(x, t)),$$

$$\Phi_x = A(\Phi),$$

$$\Phi_t = B(\Phi),$$

- Then we say that the **equation** possesses a FID.



- FIDs generalize compatible time-space decompositions requiring Hamiltonian structures, which aim to guarantee the Liouville integrability ([1, 2]).

- [1] V.I. Arnold, *Mathematical Methods of Classical Mechanics*, Springer-Verlag, New York, 1989.
- [2] W.X. Ma, in: A. Scott (Ed.), *Encyclopedia of Nonlinear Science*, Taylor & Francis, New York, 2005, pp. 450-453.

- Through FIDs, **a PDE problem can be transformed into two associated ODE problems.** Thus, the existence of solutions can be guaranteed easily by the theory of ODEs.





## Our work

- Present two classes of PDEs of specific type possessing FIDs by introducing some general ansatzes on FIDs, motivated by the works about FIDs by Ma and You et al.
- Two kinds of functions for Bäcklund transformations are taken in our constructive computation algorithm, and the associated Frobenius integrable ODEs possess **higher-order nonlinearity**.



## 2. Specific PDEs possessing FIDs

$$R(u, u_x, u_{xx}, u_{xxx}, u_{xxxx}, u_{5x}, u_{7x}, u_{9x}, u_t, u_{tt}, u_{xxt}, \dots) = 0.$$

Many interesting wave equations belong to this class of PDEs,

- ✧ *the KdV,*
- ✧ *the potential KdV,*
- ✧ *the generalized seventh-order KdV,*
- ✧ *the Burgers,*
- ✧ *the spatially periodic third-order dispersive PDE,*
- ✧ *the b-equation,*
- ...



Based on the **symmetry constrains theory** [1-4],

- [1] W.X. Ma, W. Strampp, *Physics Letters A*, 185 (1994) 277.
- [2] W.X. Ma, *Journal of the Physical Society of Japan*, 64 (1995) 1085.
- [3] W.X. Ma, Z.X. Zhou, *Journal of Mathematical Physics*, 42 (2001) 4345.
- [4] Y.B. Zeng, W.X. Ma, *Journal of Mathematical Physics*, 40 (1999) 6526.

we consider **the following case**



$$\Phi = (\phi, \psi)^T = (\phi(x, t), \psi(x, t))^T,$$

- which satisfies **two Frobenius integrable systems**:

*a real parameter*

$$\phi_x = \psi, \quad \psi_x = \lambda \phi,$$

*undetermined polynomials*

$$\phi_t = \theta_1(\phi, \psi), \quad \psi_t = \theta_2(\phi, \psi).$$

It is easy to see that the equation can be generated from the Schrödinger spectral problem with zero potential.



- Then, we have the following mixed derivatives for the considered FIDs:

$$\begin{cases} \phi_{xt} = \psi_t = \theta_2(\phi, \psi), \\ \phi_{tx} = \theta_{1,\phi}\psi + \lambda\theta_{1,\psi}\phi, \end{cases}$$

and

$$\begin{cases} \psi_{xt} = \lambda\phi_t = \lambda\theta_1(\phi, \psi), \\ \psi_{tx} = \theta_{2,\phi}\psi + \lambda\theta_{2,\psi}\phi. \end{cases}$$



- Thus, we get

$$\theta_2(\phi, \psi) = \theta_{1,\phi}\psi + \lambda\theta_{1,\psi}\phi,$$

and accordingly,

$$\lambda\theta_1(\phi, \psi) = \theta_{1,\phi\phi}\psi^2 + 2\lambda\theta_{1,\phi\psi}\phi\psi + \lambda^2\theta_{1,\psi\psi}\phi^2 + \lambda(\theta_{1,\phi}\phi + \theta_{1,\psi}\psi).$$



the only condition on  $\theta_1$



To search for  $\theta_1$  which satisfies

$$\lambda \theta_1(\phi, \psi) = \theta_{1,\phi\phi} \psi^2 + 2\lambda \theta_{1,\phi\psi} \phi\psi + \lambda^2 \theta_{1,\psi\psi} \phi^2 + \lambda(\theta_{1,\phi} \phi + \theta_{1,\psi} \psi).$$

Take the following ansatz

$$\theta_1 = \sum_{i=0}^{m_1} \sum_{j=0}^{m_2} \underbrace{b_{i,j}}_{\text{arbitrary constants}} \phi^i \psi^j,$$

natural numbers

arbitrary constants



when  $m_1 = m_2 = 7$

$$\begin{aligned}\theta_1 = & -\lambda^3 b_{1,6} \phi^7 - \lambda^3 b_{0,7} \phi^6 \psi + 3\lambda^2 b_{1,6} \phi^5 \psi^2 + 3\lambda^2 b_{0,7} \phi^4 \psi^3 - 3\lambda b_{1,6} \phi^3 \psi^4 - 3\lambda b_{0,7} \phi^2 \psi^5 \\ & + b_{1,6} \phi \psi^6 + b_{0,7} \psi^7 + \lambda^2 b_{1,4} \phi^5 + \lambda^2 b_{0,5} \phi^4 \psi - 2\lambda b_{1,4} \phi^3 \psi^2 - 2\lambda b_{0,5} \phi^2 \psi^3 + b_{1,4} \phi \psi^4 \\ & + b_{0,5} \psi^5 - \lambda b_{1,2} \phi^3 - \lambda b_{0,3} \phi^2 \psi + b_{1,2} \phi \psi^2 + b_{0,3} \psi^3 + b_{1,0} \phi + b_{0,1} \psi.\end{aligned}$$

$$\begin{aligned}\theta_2 = & \lambda(-\lambda^3 b_{0,7} \phi^6 + 6\lambda^2 b_{1,6} \phi^5 \psi + 9\lambda^2 b_{0,7} \phi^4 \psi^2 - 12\lambda b_{1,6} \phi^3 \psi^3 - 15\lambda b_{0,7} \phi^2 \psi^4 + 6b_{1,6} \phi \psi^5 \\ & + 7b_{0,7} \psi^6 + \lambda^2 b_{0,5} \phi^4 - 4\lambda b_{1,4} \phi^3 \psi - 6\lambda b_{0,5} \phi^2 \psi^2 + 4b_{1,4} \phi \psi^3 + 5b_{0,5} \psi^4 - \lambda b_{0,3} \phi^2 \\ & + 2b_{1,2} \phi \psi + 3b_{0,3} \psi^2 + b_{0,1})\phi + (-7\lambda^3 b_{1,6} \phi^6 - 6\lambda^3 b_{0,7} \phi^5 \psi + 15\lambda^2 b_{1,6} \phi^4 \psi^2 \\ & + 12\lambda^2 b_{0,7} \phi^3 \psi^3 - 9\lambda b_{1,6} \phi^2 \psi^4 - 6\lambda b_{0,7} \phi \psi^5 + b_{1,6} \psi^6 + 5\lambda^2 b_{1,4} \phi^4 \\ & + 4\lambda^2 b_{0,5} \phi^3 \psi - 6\lambda b_{1,4} \phi^2 \psi^2 - 4\lambda b_{0,5} \phi \psi^3 + b_{1,4} \psi^4 \\ & - 3\lambda b_{1,2} \phi^2 - 2\lambda b_{0,3} \phi \psi + b_{1,2} \psi^2 + b_{1,0})\psi.\end{aligned}$$





- On the other hand, we consider the following specific type of the polynomial  $R$

$$\begin{aligned} R(u, u_x, u_{xx}, u_{xxx}, u_{xxxx}, u_{5x}, u_{7x}, u_{9x}, u_t, u_{tt}, u_{xxt}, \dots) \\ = d_{0,0}u + d_{0,1}u_x + d_{0,2}u_{xx} + d_{0,3}u_{xxx} + d_{0,4}u_{xxxx} + d_{1,0}u^2 + d_{1,1}uu_x + d_{1,2}uu_{xx} + d_{1,3}uu_{xxx} \\ + d_{1,4}uu_{xxxx} + d_{2,0}u^2_x + d_{2,1}u_xu_{xx} + d_{2,2}u_xu_{xxx} + d_{2,3}u_xu_{xxxx} + d_{3,0}u^2_{xx} + d_{3,1}u_{xx}u_{xxx} \\ + d_{3,2}u_{xx}u_{xxxx} + d_{4,0}u^2_{xxx} + d_{4,1}u_{xxx}u_{xxxx} + d_{5,0}u^2_{xxxx} + e_0u_{5x} + e_1u_{7x} + e_2u_{9x} \\ + e_3uu_{5x} + e_4uu_{7x} + e_5uu_{9x} + f_0u^2u_x + f_1u^2u_{xxx} + f_2u^2u_{5x} + f_3u^3_x \\ + f_4u^3u_x + f_5uu_xu_{xx} + g_0u_t + g_1u_{tt} + g_2u_{xxt} + g_3u^2_{xxt}. \end{aligned}$$



- **Theorem (a)** If we take a Bäcklund transformation

$$u = (\ln \phi)_x = \frac{\psi}{\phi}$$

from the Frobenius integrable systems

$$\phi_x = \psi, \quad \psi_x = \lambda \phi,$$

$$\phi_t = \theta_1(\phi, \psi), \quad \psi_t = \theta_2(\phi, \psi),$$

to the PDE

$$P(u, u_t, u_x, u_{xt}, \dots) = 0,$$



$$\begin{aligned}\theta_1 = & -\lambda^3 b_{1,6} \phi^7 + 3\lambda^2 b_{1,6} \phi^5 \psi^2 - 3\lambda b_{1,6} \phi^3 \psi^4 + b_{1,6} \phi \psi^6 + \lambda^2 b_{1,4} \phi^5 - 2\lambda b_{1,4} \phi^3 \psi^2 \\ & + b_{1,4} \phi \psi^4 - \lambda b_{1,2} \phi^3 + b_{1,2} \phi \psi^2 + b_{1,0} \phi + b_{0,1} \psi ,\end{aligned}$$

$$\begin{aligned}\theta_2 = & \lambda(6\lambda^2 b_{1,6} \phi^5 \psi - 12\lambda b_{1,6} \phi^3 \psi^3 + 6b_{1,6} \phi \psi^5 - 4\lambda b_{1,4} \phi^3 \psi + 4b_{1,4} \phi \psi^3 + 2b_{1,2} \phi \psi + b_{0,1}) \phi \\ & + (15\lambda^2 b_{1,6} \phi^4 \psi^2 - 7\lambda^3 b_{1,6} \phi^6 - 9\lambda b_{1,6} \phi^2 \psi^4 + b_{1,6} \psi^6 + 5\lambda^2 b_{1,4} \phi^4 - 6\lambda b_{1,4} \phi^2 \psi^2 \\ & + b_{1,4} \psi^4 - 3\lambda b_{1,2} \phi^2 + b_{1,2} \psi^2 + b_{1,0}) \psi ,\end{aligned}$$



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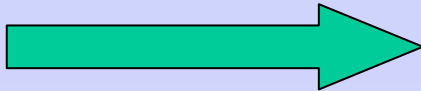


$$\begin{aligned} R = & (224\lambda^2 e_4 + 8\lambda d_{2,3} - 40\lambda e_3 + 24d_{0,4} - 2d_{2,1} - 6d_{1,3})u^3 u_x + f_1 u^2 u_{xxx} + f_2 u^2 u_{5x} + f_3 u_x^3 \\ & + d_{1,4} u u_{xxxx} + d_{1,3} u u_{xxx} + (3d_{2,2} - 40\lambda f_2 - 60e_0 + 2d_{3,0} - 3f_1 - 1/2 f_3 + 12d_{1,4}) u u_x u_{xx} \\ & + d_{1,2} u u_{xx} + e_3 u u_{5x} + e_4 u u_{7x} + 630e_2 u_{xxxx}^2 - 35e_4 u_{xxx} u_{xxxx} + d_{2,3} u_x u_{xxxx} + d_{4,0} u_{xxx}^2 \\ & + d_{2,2} u_x u_{xxx} + d_{3,0} u_{xx}^2 + d_{2,1} u_x u_{xx} - (256\lambda^4 e_2 + 64\lambda^3 e_1 + 16\lambda^3 f_2 + 16\lambda^2 e_0 \\ & + 4\lambda^2 f_1 - 8\lambda^2 d_{1,4} + 4\lambda b_{0,1} g_2 + 4\lambda d_{0,3} - 2\lambda d_{1,2} + b_{0,1} g_0 + d_{0,1}) u^2 u_x / \lambda \\ & + (24\lambda^2 e_3 - 288\lambda^3 e_4 - 8\lambda^2 d_{2,3} + 2b_{0,1}^2 g_1 + 2\lambda d_{2,1} - 16\lambda d_{0,4} + 2\lambda d_{1,3} \\ & + 2d_{0,2}) u u_x + (105e_1 - 3/4 d_{4,0} + 5/2 f_2 - 3/4 b_{0,1}^2 g_3) u_{xx} u_{xxxx} \\ & + d_{0,4} u_{xxxx} - (140\lambda e_4 + 2d_{2,3} + 10e_3) u_{xx} u_{xxx} + d_{0,3} u_{xxx} \\ & + d_{0,2} u_{xx} + e_2 u_{9x} + e_1 u_{7x} + e_0 u_{5x} + d_{0,1} u_x - (4\lambda^3 c_{0,1}^2 g_3 \\ & + 7936\lambda^4 e_2 + 4\lambda^3 d_{4,0} - 272\lambda^3 e_1 + 16\lambda^2 e_0 \\ & - 2\lambda^2 d_{2,2} + \lambda^2 f_3 - 2\lambda b_{0,1} g_2 - 2\lambda d_{0,3} \\ & + b_{0,1} g_0 + d_{0,1}) u_x^2 / \lambda + g_0 u_t + g_1 u_{tt} \\ & + g_2 u_{xxt} + g_3 u_{xxt}^2. \end{aligned}$$



- **(b)** If we take a Bäcklund transformation

$$u = (\ln \phi)_{xx} = \frac{\lambda \phi^2 - \psi^2}{\phi^2}$$



$$\theta_1 = -\lambda^3 b_{1,6} \phi^7 + 3\lambda^2 b_{1,6} \phi^5 \psi^2 - 3\lambda b_{1,6} \phi^3 \psi^4 + b_{1,6} \phi \psi^6 + \lambda^2 b_{1,4} \phi^5 - 2\lambda b_{1,4} \phi^3 \psi^2 + b_{1,4} \phi \psi^4 - \lambda b_{1,2} \phi^3 + b_{1,2} \phi \psi^2 + b_{1,0} \phi + b_{0,1} \psi,$$

$$\theta_2 = -\lambda^3 b_{1,6} \phi^6 \psi + 3\lambda^2 b_{1,6} \phi^4 \psi^3 - 3\lambda b_{1,6} \phi^2 \psi^5 + \lambda^2 b_{1,4} \phi^4 \psi - 2\lambda b_{1,4} \phi^2 \psi^3 - \lambda b_{1,2} \phi^2 \psi + \lambda b_{0,1} \phi + b_{1,6} \psi^7 + b_{1,4} \psi^5 + b_{1,2} \psi^3 + b_{0,1} \psi,$$



$R =$

$$\begin{aligned} & f_4 u^3 u_x + f_1 u^2 u_{xxx} + f_5 u u_x u_{xx} + (443520 \lambda^2 e_2 - 352 \lambda^2 d_{4,1} - 3584 \lambda^2 e_4 - 120 \lambda e_3 + 8 \lambda f_1 - \lambda f_4 + 2 \lambda f_5 \\ & - 360 e_0 + 12 d_{1,3} + 6 d_{2,1}) u^2 u_x + (4 d_{4,1} - 5040 e_2 + 56 e_4) u^2 u_{5x} + e_4 u u_{7x} + e_3 u u_{5x} + d_{1,4} u u_{xxxx} + d_{1,3} u u_{xxx} \\ & + d_{1,2} u u_{xx} + (240 \lambda d_{4,1} - 302400 \lambda e_2 + 1680 \lambda e_4 + 30 d_{2,3} + 18 d_{3,1} - 5040 e_1 + 90 e_3 - 3 f_1 + 1/4 f_4 \\ & - 3/2 f_5) u_x^3 + d_{2,3} u_x u_{xxxx} + d_{2,1} u_x u_{xx} + d_{3,2} u_{xx} u_{xxxx} + d_{4,1} u_{xxx} u_{xxxx} + d_{3,0} u_{xx}^2 + d_{3,1} u_{xx} u_{xxx} + (32 \lambda^3 d_{3,2} \\ & - 16 \lambda^2 d_{1,4} - 4 \lambda^2 d_{3,0} + 6 b_{0,1}^2 g_1 + 2 \lambda d_{1,2} + 6 d_{0,2}) u^2 + (65280 \lambda^3 e_2 - 64 \lambda^3 d_{4,1} - 64 \lambda^3 e_4 \\ & - 16 \lambda^2 d_{2,3} - 16 \lambda^2 d_{3,1} + 4032 \lambda^2 e_1 - 16 \lambda^2 e_3 + 240 \lambda e_0 - 4 \lambda d_{1,3} - 4 \lambda d_{2,1} + 12 b_{0,1} g_2 \\ & + 12 d_{0,3}) u u_x + (10 \lambda d_{1,4} - 4 \lambda^2 d_{3,2} + 30 d_{0,4} - 3/2 d_{1,2}) u_x^2 - (5/2 d_{1,4} + 3/4 d_{3,0}) \\ & u_x u_{xxx} + e_2 u_{9x} + e_1 u_{7x} + e_0 u_{5x} + d_{0,4} u_{xxxx} + d_{0,3} u_{xxx} + d_{0,2} u_{xx} - (b_{0,1}^2 g_3 + 5/4 d_{3,2}) \\ & u_{xxx}^2 - (4 \lambda b_{0,1}^2 g_1 + 16 \lambda^2 d_{0,4} + 4 \lambda d_{0,2}) u - (256 \lambda^4 e_2 + 64 \lambda^3 e_1 + 16 \lambda^2 e_0 \\ & + 4 \lambda b_{0,1} g_2 + 4 \lambda d_{0,3} + b_{0,1} g_0) u_x + g_0 u_t + g_1 u_{tt} + g_2 u_{xxt} + g_3 u_{xxt}^2. \end{aligned}$$



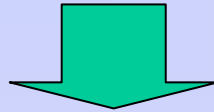
- Take **a special reduction** as follows:

$$\begin{aligned} & R(u, u_x, u_{xx}, u_{xxx}, u_{xxxx}, u_{5x}, u_{7x}, u_{9x}, u_t, u_{tt}, u_{xxt}, \dots) \\ &= d_{0,0}u + d_{0,1}u_x + d_{0,2}u_{xx} + d_{0,3}u_{xxx} + d_{0,4}u_{xxxx} + d_{1,1}uu_x + d_{1,2}uu_{xx} + d_{1,3}uu_{xxx} \\ &+ d_{2,0}u_x^2 + d_{2,1}u_xu_{xx} + d_{2,2}u_xu_{xxx} + d_{2,3}u_xu_{xxxx} + d_{3,1}u_{xx}u_{xxx} + e_0u_{5x} + e_1u_{7x} \\ &+ e_2u_{9x} + e_3uu_{5x} + f_0u^2u_x + f_1u^2u_{xxx} + f_2u^2u_{5x} + f_3u_x^3 + f_4u^3u_x \\ &+ f_5uu_xu_{xx} + g_0u_t + g_1u_{tt} + g_2u_{xxt}. \end{aligned}$$



For the  $u = (\ln \phi)_x = \frac{\psi}{\phi}$

**The corresponding results**



$$\theta_1 = -\lambda^3 b_{1,6} \phi^7 + 3 \lambda^2 b_{1,6} \phi^5 \psi^2 - 3 \lambda b_{1,6} \phi^3 \psi^4 + b_{1,6} \phi \psi^6 + \lambda^2 b_{1,4} \phi^5 \\ - 2 \lambda b_{1,4} \phi^3 \psi^2 + b_{1,4} \phi \psi^4 - \lambda b_{1,2} \phi^3 + b_{1,2} \phi \psi^2 + b_{1,0} \phi + b_{0,1} \psi,$$

$$\theta_2 = \lambda (6 \lambda^2 b_{1,6} \phi^5 \psi - 12 \lambda b_{1,6} \phi^3 \psi^3 + 6 b_{1,6} \phi \psi^5 - 4 \lambda b_{1,4} \phi^3 \psi + 4 b_{1,4} \phi \psi^3 + 2 b_{1,2} \phi \psi \\ + b_{0,1}) \phi + (-7 \lambda^3 b_{1,6} \phi^6 + 15 \lambda^2 b_{1,6} \phi^4 \psi^2 - 9 \lambda b_{1,6} \phi^2 \psi^4 + b_{1,6} \psi^6 + 5 \lambda^2 b_{1,4} \phi^4 \\ - 6 \lambda b_{1,4} \phi^2 \psi^2 + b_{1,4} \psi^4 - 3 \lambda b_{1,2} \phi^2 + b_{1,2} \psi^2 + b_{1,0}) \psi,$$





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$$\begin{aligned} R = & (8 \lambda d_{2,3} - 40 \lambda e_3 + 24 d_{0,4} - 6 d_{1,3} - 2 d_{2,1}) u^3 u_x + f_1 u^2 u_{xxx} - 42 e_1 u^2 u_{5x} + f_3 u_x^3 \\ & + (336 \lambda^2 e_1 - 2 \lambda d_{2,2} - 4 \lambda f_1 + \lambda f_3 - 6 b_{0,1} g_2 - 6 d_{0,3} + 2 d_{1,2} + d_{2,0}) u^2 u_x \\ & + d_{1,2} u u_{xx} + (1680 \lambda e_1 + 3 d_{2,2} - 60 e_0 - 3 f_1 - 1/2 f_3) u u_x u_{xx} + d_{1,3} u u_{xxx} \\ & + e_3 u u_{5x} + d_{2,1} u_x u_{xx} + d_{2,2} u_x u_{xxx} + d_{2,3} u_x u_{xxxx} + d_{2,0} u_x^2 + (24 \lambda^2 e_3 \\ & - 8 \lambda^2 d_{2,3} + 2 b_{0,1}^2 g_1 - 16 \lambda d_{0,4} + 2 \lambda d_{1,3} + 2 \lambda d_{2,1} + 2 d_{0,2}) u u_x \\ & - (2 d_{2,3} + 10 e_3) u_{xx} u_{xxx} + d_{0,2} u_{xx} + d_{0,3} u_{xxx} + d_{0,4} u_{xxxx} \\ & + e_0 u_{5x} + e_1 u_{7x} + (272 \lambda^3 e_1 + 2 \lambda^2 d_{2,2} - 16 \lambda^2 e_0 \\ & - \lambda^2 f_3 + 2 \lambda b_{0,1} g_2 + 2 \lambda d_{0,3} - \lambda d_{2,0} \\ & - b_{0,1} g_0) u_x + g_0 u_t + g_1 u_{tt} \\ & + g_2 u_{xxt}. \end{aligned}$$



- It includes many nonlinear equations,
- **the potential KdV equation**

$$g_0 u_t + d_{0,3} u_{xxx} + 6d_{0,3} u_x^2 = 0,$$

- **the Burgers equation**

$$g_0 u_t + 2d_{0,2} uu_x + d_{0,2} u_{xx} = 0,$$

- **b - equation**

$$g_0 u_t - \frac{1}{2\lambda} u_{xxt} - 2uu_x = \frac{1}{2\lambda} (-3u_x u_{xx} + uu_{xxx}).$$

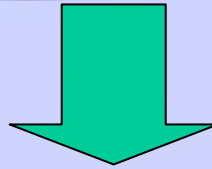
a general case of the famous Rod equation [1]

[1] H.H. Dai, Y. Huo, *Proceedings of the Royal Society of London Series A*, 456 (2000) 331.



For the  $u = (\ln \phi)_{xx} = \frac{\lambda \phi^2 - \psi^2}{\phi^2}$

**The corresponding results**



$$\theta_1 = -\lambda^3 b_{1,6} \phi^7 + 3 \lambda^2 b_{1,6} \phi^5 \psi^2 - 3 \lambda b_{1,6} \phi^3 \psi^4 + b_{1,6} \phi \psi^6 + \lambda^2 b_{1,4} \phi^5 \\ - 2 \lambda b_{1,4} \phi^3 \psi^2 + b_{1,4} \phi \psi^4 - \lambda b_{1,2} \phi^3 + b_{1,2} \phi \psi^2 + b_{1,0} \phi + b_{0,1} \psi,$$

$$\theta_2 = -\lambda^3 b_{1,6} \phi^6 \psi + 3 \lambda^2 b_{1,6} \phi^4 \psi^3 - 3 \lambda b_{1,6} \phi^2 \psi^5 + \lambda^2 b_{1,4} \phi^4 \psi - 2 \lambda b_{1,4} \phi^2 \psi^3 \\ - \lambda b_{1,2} \phi^2 \psi + \lambda b_{0,1} \phi + b_{1,6} \psi^7 + b_{1,4} \psi^5 + b_{1,2} \psi^3 + \psi b_{1,0},$$



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$$\begin{aligned} R = & f_4 u^3 u_x + f_1 u^2 u_{xxx} + (443520 \lambda^2 e_2 - 120 \lambda e_3 + 8 \lambda f_1 - \lambda f_4 + 2 \lambda f_5 + 12 d_{1,3} + 6 d_{2,1} \\ & - 360 e_0) u^2 u_x + d_{1,2} u u_{xx} + d_{1,3} u u_{xxx} + e_3 u u_{5x} + d_{3,0} u_{xx} u_{xxx} + d_{2,1} u_x u_{xx} + (30 d_{2,3} \\ & + 18 d_{3,0} - 302400 \lambda e_2 - 5040 e_1 + 90 e_3 - 3 f_1 + 1/4 f_4 - 3/2 f_5) u_x^3 + e_2 u_{9x} \\ & + d_{0,3} u_{xxx} + f_5 u u_x u_{xx} + d_{0,4} u_{xxxx} + e_0 u_{5x} + e_1 u_{7x} + (30 d_{0,4} - 3/2 d_{1,2}) u_x^2 \\ & + (4/3 \lambda^2 d_{1,2} - 16 \lambda^2 d_{0,4}) u - (b_{0,1}^2 g_1 + 1/3 \lambda d_{1,2}) u_{xx} - (64 \lambda^3 e_1 \\ & + 256 \lambda^4 e_2 + 16 \lambda^2 e_0 + 4 \lambda c_{0,1} g_2 + 4 \lambda d_{0,3} + b_{0,1} g_0) u_x \\ & + d_{2,3} u_x u_{xxxx} - 5040 e_2 u^2 u_{5x} + (65280 \lambda^3 e_2 \\ & - 16 \lambda^2 d_{2,3} - 16 \lambda^2 d_{3,0} + 4032 \lambda^2 e_1 \\ & - 16 \lambda^2 e_3 - 4 \lambda d_{1,3} - 4 \lambda d_{2,1} \\ & + 240 \lambda e_0 + 12 b_{0,1} g_2 \\ & + 12 d_{0,3}) u u_x + g_0 u_t \\ & + g_1 u_{tt} + g_2 u_{xxt}. \end{aligned}$$



It includes:

- **the KdV equation**

$$g_0 u_t + 12d_{0,3} u u_x + d_{0,3} u_{xxx} = 0,$$

- **the family of spatially periodic third-order dispersive PDE**

$$b_{0,1} g_0 u_t - b_{0,1}^2 g_0 u_x + b_{0,1} d_{0,3} u_{xxx} - d_{0,3} u_{xxt} = b_{0,1} \alpha \left( \lambda u^2 + u_x^2 - \frac{1}{2} u u_{xx} \right)_x,$$



- the generalized seventh-order KdV equation

$$g_0 u_t + (2\lambda f_5 - 120\lambda e_3 + 8\lambda f_1)u^2 u_x + (30d_{2,3} + 18d_{3,0} - \frac{3}{2}f_5 + 90e_3 - 5040e_1 - 3f_1)u_x^3 + f_5 u u_x u_{xx} + f_1 u^2 u_{xxx} + d_{3,0} u_{xx} u_{xxx} + d_{2,3} u_x u_{xxxx} + e_3 u u_{5x} + e_1 u_{7x} = 0,$$

which has two well-known special cases [1]:

→ **the Lax seventh-order equation**

→ **the Sawada-Kotera-Ito seventh-order equation**

[1] M. Ito, *Journal of the Physical Society of Japan*, 49 (1980) 771



## 3. Conclusions

- Through two Bäcklund transformations of dependent variables, ninth-order PDEs possessing the FIDs have been obtained.
- Two special classes of such nonlinear PDEs possessing special FIDs have been presented under a reduction.
- The obtained PDEs contain various significant nonlinear wave equations.



- The approach adopted here can be easily applied to nonlinear variable coefficient PDEs, which are quite intriguing in *ocean dynamics, fluid mechanics, plasma physics, etc.*

- **A question ?** *If we take the Möbius transformation*

$$u = (a\varphi + b\psi) / (c\varphi + d\psi), \quad (ad - bc \neq 0)$$

instead of the logarithmic derivative type Bäcklund transformations, what we will get?





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- *Thank Prof. Wen-Xiu Ma, Prof. Xingbiao Hu and Prof. Qingping Liu!*
- **Thank you for your attention!**