

# DARBOUX TRANSFORMATION FOR THE NONLINEAR SCHRÖDINGER EQUATION

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(joint with C. van der Mee, Cagliari)

**Outline:**

- Darboux transformation
- Bäcklund transformation
- Darboux transform for Zakharov-Shabat system/NLS equation
- general approach to Darboux transformations
- references

**For details:**

T. Aktosun and C. van der Mee, *A unified approach to Darboux transformations*, Inverse Problems (2009), to appear

**Unperturbed problem:**  $\mathcal{L}\psi = \lambda\psi$

- $\mathcal{L}$  differential operator in  $x$
- $\lambda$  spectral parameter
- potential  $u(x, t)$ ,  $t$  parameter
- $\psi(\lambda, x, t)$  wave function
- spectrum of  $\mathcal{L}$ , discrete and continuous spectra
- eigenvalues and eigenfunctions
- (generalized) eigenvalues and eigenfunctions

**Perturbed problem:**  $\tilde{\mathcal{L}}\tilde{\psi} = \lambda\tilde{\psi}$

- $\tilde{\mathcal{L}}$  is a finite-rank perturbation of  $\mathcal{L}$
- $\tilde{\mathcal{L}}$  is obtained from  $\mathcal{L}$  by changing only the discrete spectrum
- perturbed potential  $\tilde{u}(x, t)$
- perturbed  $\tilde{\psi}(\lambda, x, t)$  wave function

## **Darboux transformation**

- express, via unperturbed quantities and finite-rank perturbation:

$$\tilde{u}(x, t) - u(x, t) \quad (\text{DT at the potential level})$$

$$\tilde{\psi}(\lambda, x, t) - \psi(\lambda, x, t) \quad (\text{DT at the wave function level})$$

## **Bäcklund transformation**

- differential equation(s) only involving:

$$\tilde{u}(x, t), u(x, t), \text{ and their } x\text{- and } t\text{-derivatives}$$

**Zakharov-Shabat system:**  $\mathcal{L}\psi = \lambda\psi$

- $\mathcal{L} := i \begin{bmatrix} \frac{d}{dx} & -u(x, t) \\ -u(x, t)^* & -\frac{d}{dx} \end{bmatrix}$

- $\psi(\lambda, x, t) := \begin{bmatrix} \psi_1(\lambda, x, t) \\ \psi_2(\lambda, x, t) \end{bmatrix}$  Jost solution (from the left)

- $\psi(\lambda, x, t) \sim \begin{bmatrix} 0 \\ e^{i\lambda x} \end{bmatrix}$  as  $x \rightarrow +\infty$

**Nonlinear Schrödinger (NLS) equation:**  $iu_t + u_{xx} + 2|u|^2u = 0$

## Inverse scattering transform

- integrable nonlinear PDE, linear spectral problem  $\mathcal{L}\psi = \lambda\psi$

- KdV  $u_t - 6uu_x + u_{xxx} = 0$

1-D Schrödinger eq  $-\frac{d^2\psi}{dx^2} + u(x,t)\psi = \lambda\psi$

- mKdV  $u_t + 6u^2u_x + u_{xxx} = 0,$  
$$\begin{cases} \frac{d\xi}{dx} = -i\lambda\xi + u(x,t)\eta \\ \frac{d\eta}{dx} = i\lambda\eta - u(x,t)\xi \end{cases}$$

- sine-Gordon  $u_{xt} = \sin u,$  
$$\begin{cases} \frac{d\xi}{dx} = -i\lambda\xi - \frac{1}{2}u_x(x,t)\eta \\ \frac{d\eta}{dx} = i\lambda\eta + \frac{1}{2}u_x(x,t)\xi \end{cases}$$

## Darboux transformation for ZS system/NLS

Add one bound state at  $\lambda = \lambda_1$  with the norming constant  $c_1(t)$ . Then:

$$\tilde{u}(x, t) - u(x, t) = \frac{P_0}{|\Gamma_1|^2 + |c_1|^2 \Gamma_2^2}$$

$$\begin{bmatrix} \tilde{\psi}_1(\lambda, x, t) \\ \tilde{\psi}_2(\lambda, x, t) \end{bmatrix} - \begin{bmatrix} \psi_1(\lambda, x, t) \\ \psi_2(\lambda, x, t) \end{bmatrix} = \frac{1}{|\Gamma_1|^2 + |c_1|^2 \Gamma_2^2} \begin{bmatrix} P_1 & P_2 \\ P_3 & P_4 \end{bmatrix} \begin{bmatrix} P_5 \\ P_6 \end{bmatrix}$$

$$\begin{aligned}
\Gamma_1 &:= 1 + ic_1 \left[ \psi_1(\lambda_1, x) \dot{\psi}_2(\lambda_1, x) - \psi_2(\lambda_1, x) \dot{\psi}_1(\lambda_1, x) \right] \\
\Gamma_2 &:= \frac{|\psi_1(\lambda_1, x)|^2 + |\psi_2(\lambda_1, x)|^2}{2 \operatorname{Im}[\lambda_1]} \\
P_0 &:= 2c_1 \psi_1(\lambda_1, x)^2 \Gamma_1^* - 2c_1^* [\psi_2(\lambda_1, x)^*]^2 \Gamma_1 + 4|c_1|^2 \psi_1(\lambda_1, x) \psi_2(\lambda_1, x)^* \Gamma_2 \\
P_1 &:= -|c_1|^2 \psi_2(\lambda_1, x)^* \Gamma_2 - c_1 \psi_1(\lambda_1, x) \Gamma_1^* \\
P_2 &:= |c_1|^2 \psi_1(\lambda_1, x) \Gamma_2(x) - c_1^* \psi_2(\lambda_1, x)^* \Gamma_1 \\
P_3 &:= |c_1|^2 \psi_1(\lambda_1, x)^* \Gamma_2(x) - c_1 \psi_2(\lambda_1, x) \Gamma_1^* \\
P_4 &:= |c_1|^2 \psi_2(\lambda_1, x) \Gamma_2(x) + c_1^* \psi_1(\lambda_1, x)^* \Gamma_1 \\
P_5 &:= \frac{i}{\lambda - \lambda_1} [\psi_1(\lambda_1, x) \psi_2(\lambda, x) - \psi_2(\lambda_1, x) \psi_1(\lambda, x)] \\
P_6 &:= \frac{-i}{\lambda - \lambda_1^*} [\psi_1(\lambda_1, x)^* \psi_1(\lambda, x) + \psi_2(\lambda_1, x)^* \psi_2(\lambda, x)] \\
c_1 &:= c_1(t) = c_1(0) e^{4i\lambda_1^2 t} \\
\psi_1(\lambda_1, x) &:= \psi_1(\lambda_1, x, t), \quad \psi_2(\lambda, x) := \psi_2(\lambda, x, t)
\end{aligned}$$

## Darboux transformation for ZS system/NLS

- Add/remove  $N$  bound states at  $\{\lambda_j\}_{j=1}^N$  with  $\lambda_j$  of multiplicity  $n_j$  and norming constants  $c_{j1}, c_{j2}, \dots, c_{jn_j}$
- determine  $\tilde{u}(x, t) - u(x, t)$  and  $\tilde{\psi}(\lambda, x, t) - \psi(\lambda, x, t)$

## General approach to Darboux transformations

- unified approach
- applicable when Marchenko/GL methods are applicable
- multiple eigenvalues with multiplicities
- applicable to matrix versions of integrable equations
- no assumptions on extensions to both  $\mathbf{C}^+$  and  $\mathbf{C}^-$

## General approach to Darboux transformations

- relate  $\psi(\lambda, x, t)$  to  $\alpha(x, y, t)$  via Fourier transform
- $\tilde{\psi}(\lambda, x, t)$  is related to  $\tilde{\alpha}(x, y, t)$
- $\alpha(x, y, t)$  satisfies  $\alpha + \omega + \alpha\Omega = 0$  (Marchenko/Gel'fand-Levitan eq)
 
$$\alpha(x, y, t) + \omega(x, y, t) + \int_x^\infty dz \alpha(x, z, t) \omega(z, y, t) = 0$$
- $\tilde{\alpha}(x, y, t)$  satisfies  $\tilde{\alpha} + \tilde{\omega} + \tilde{\alpha}\tilde{\Omega} = 0$ 

$$\tilde{\Omega} - \Omega = FG, \quad \tilde{\omega}(x, y, t) - \omega(x, y, t) = f(x, t) g(y, t)$$
- $\tilde{\alpha}[I + FG(I + R)] = \alpha - fg(I + R)$  equivalent separable eq
 
$$I + R := (I + \Omega)^{-1}$$
 explicit evaluation in terms of  $\alpha$
- $\tilde{\alpha}(x, y, t) - \alpha(x, y, t)$  is obtained in terms of  $\alpha(x, y, t)$ ,  $f(x, t)$ ,  $g(y, t)$
- $\tilde{\psi}(\lambda, x, t) - \psi(\lambda, x, t)$  is obtained from  $\tilde{\alpha}(x, y, t) - \alpha(x, y, t)$
- $\tilde{u}(x, t) - u(x, t)$  is obtained from  $\tilde{\alpha}(x, x, t) - \alpha(x, x, t)$

## General approach to Darboux transformations

- $\alpha + \omega + \alpha\Omega = 0 \quad \alpha = -\omega(I + \Omega)^{-1} = -\omega(I + R)$

- $R + \Omega + R\Omega = 0$

$$r(x; y, z) + \omega(y, z) + \int_x^\infty ds r(x; y, s) \omega(s, z) = 0, \quad x < \min\{y, z\}$$

$$r(x; y, z) = \begin{cases} \alpha(y, z) + \int_x^y ds J \alpha(s, y)^\dagger J \alpha(s, z), & x < y < z, \\ J \alpha(z, y)^\dagger J + \int_x^z ds J \alpha(s, y)^\dagger J \alpha(s, z), & x < z < y, \end{cases}$$

- $\tilde{\alpha} + \tilde{\omega} + \tilde{\alpha}\tilde{\Omega} = 0 \quad \tilde{\alpha} + \omega + fg + \tilde{\alpha}(\Omega + FG) = 0$

$$[\tilde{\alpha} + \omega + fg + \tilde{\alpha}(\Omega + FG)](I + R) = 0$$

$$\tilde{\alpha}[I + FG(I + R)] = \alpha - fg(I + R) \quad \text{separable integral eq}$$

$$\tilde{\alpha} - \alpha = -(f + \alpha F)[I + G(I + R)F]^{-1}G(I + R)$$

## General approach applied to ZS/NLS

$$\bullet f(x, t) = \begin{bmatrix} 0 & B^\dagger e^{-A^\dagger x} \\ C e^{-Ax - 4iA^2 t} & 0 \end{bmatrix}$$

$$\bullet g(y, t) = \begin{bmatrix} e^{-Ay} B & 0 \\ 0 & -e^{-A^\dagger y + 4i(A^\dagger)^2 t} C^\dagger \end{bmatrix}$$

$$\bullet \alpha(x, y, t) = [\bar{K}(x, y, t) \quad K(x, y, t)] \quad 2 \times 2 \text{ matrix}$$

$$\bullet \omega(x, y, t) = \begin{bmatrix} 0 & -\Omega_1(x + y, t)^\dagger \\ \Omega_1(x + y, t) & 0 \end{bmatrix} \quad 2 \times 2 \text{ matrix}$$

$$\bullet u(x, t) = -2 [1 \quad 0] \alpha(x, x, t) \begin{bmatrix} 0 \\ 1 \end{bmatrix}$$

$$\bullet \begin{cases} \bar{K}(x, y, t) + \begin{bmatrix} 0 \\ \Omega_1(x + y, t) \end{bmatrix} + \int_x^\infty dz K(x, z, t) \Omega_1(z + y, t) = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \\ K(x, y, t) - \begin{bmatrix} \Omega_1(x + y, t)^\dagger \\ 0 \end{bmatrix} - \int_x^\infty dz \bar{K}(x, z, t) \Omega_1(z + y, t)^\dagger = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \end{cases}$$

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