Mineralogy of a natural As-rich hydrous ferric oxide coprecipitate formed by mixing of hydrothermal fluid and seawater: Implications regarding surface complexation and color banding in ferrihydrite deposits

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ABSTRACT

We characterized the most As-rich natural hydrous ferric oxide (HFO) material ever reported using powder X-ray diffraction (pXRD), transmission electron microscopy (TEM), X-ray fluorescence spectroscopy (XRF), light element analysis using gas chromatography (GC), visible-infrared (vis-IR) diffuse reflectivity, 57Fe Mössbauer spectroscopy, and superconducting quantum interference device (SQUID) magnetometry. We find that the natural As-HFO material is very similar to synthetic coprecipitated As-HFO materials, but is significantly different from all known natural and synthetic As-free HFO materials and ferrihydrite samples. The pXRD patterns show systematic differences with patterns for 2-line ferrihydrite, that are interpreted as evidence for significant populations of oxygen-coordinated Fe-As pairs. Observations by TEM, combined with energy dispersive spectroscopy (EDS) microanalysis, show agglomerations of nanophase primary particles and no evidence for other Fe- or As-bearing phases. Mössbauer spectroscopy shows octahedrally coordinated Fe³⁺, with a large fraction (~20%) of the octahedral Fe environments that are significantly distorted by the presence of As, compared to the Fe local environments in As-free ferrihydrite and HFO samples. The loss on ignition (LOI) is quantitatively consistent with OH + H₂O, measured by GC, which, in turn, is consistent with ~1 nm diameter primary particles having all their surface cations (Fe³⁺, As⁵⁺, Si⁴⁺, C⁴⁺) coordinated on the free surface side by OH⁻ and OH₂⁻. The banding into adjacent yellowish and reddish layers that occurs in the As-HFO deposits was studied by performing mineralogical analyses of the separated adjacent layers of two couplets of yellowish and reddish material. The yellowish samples were found not to contain secondary crystalline phases (as did the reddish samples, in small amounts) and to be relatively As-rich, C- and Si-poor. The observed anticorrelations between As and Si and between As and inorganic C suggest that natural HFOs, which usually contain significant molar amounts of Si, may not be as efficient at surface complexing As (and P) as their Si and C-free synthetic counterparts, unless formed by co-precipitation with the As (or P). The yellowish and reddish layers were also clearly resolved by both Mössbauer spectroscopy and magnetometry. Complexation of arsenate onto the HFO core was found to significantly increase the average quadrupole splitting (QS) obtained from Mössbauer spectroscopy by an amount that could not be explained by other chemical differences and that is consistent with an ~1 nm diameter particle size and somewhat smaller HFO core. The Munsell hue YR index (5–10 YR) was found to be strongly correlated to the average QS, thereby establishing that the color differences, corresponding to the measured shifts of the main visible band edge, are due to the local distortions in the ¹⁰⁶Fe³⁺ environments that are induced by As complexation, via their influence on the relevant ligand field transitions. SQUID magnetometry allows the following observations. (1) The superparamagnetic to superferromagnetic transitions occur at 25 K and lower in As-HFO, compared to 55 K in synthetic 2-line ferrihydrite, suggesting a smaller magnetic primary particle (or core) size for As-HFO and inter-particle magnetic interaction reduction by surface complexed As, Si, and C. (2) The ratio of supermoment magnitude to magnetic particle size (m²/n, where m is the net number of Fe³⁺ atomic moments per supermoment and n is the number of Fe³⁺ cations per particle or HFO core) decreases with increasing As content in the sequence synthetic-HFO > reddish-As-HFO > yellowish-As-HFO. (3) The magnetic susceptibility magnitudes for As-HFO and synthetic 2-line ferrihydrite differ by a factor of 10 and suggest different supermoment formation mechanisms (m²/n < 1 vs. m²/n > 1, respectively) related to differences in intra-particle cationic and anionic disorder and magnetic particle size.

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INTRODUCTION

The mobility and immobilization of As in natural waters is, in most circumstances, controlled by sorption onto particulate phases (Mok and Wai 1994). In particular, hydrous ferric oxide (HFO) is known to have a great sorption affinity for both arsenite [As(III)] and arsenate [As(V)] and sorption onto HFO is believed to be a dominant mechanism controlling As cycling in many natural environments (Cullen and Reimer 1989; Maher and Butler 1988). Dominant As/HFO associations have been documented (or inferred) to occur in: lacustrine sediments (e.g., Aggett and O’Brian 1985; Belzile and Tessier 1990; Peterson and Carpenter 1986), marine sediments (e.g., Belzile 1988; Edenborn et al. 1986), soils (e.g., Hess and Blanchard 1976; Bowell 1994) and acid mine drainage environments (e.g., Leblanc et al. 1996). In addition, sorption onto HFO is believed to be the main mechanism in the common ferric chloride coagulant-based water treatment method for As removal (Hering et al. 1996, 1997) and the most effective As removal treatment of spent hydrothermal fluid in hydrothermal power generation plants has been found to be As/HFO coprecipitation combined with a standard flotation separation (De Carlo and Thomas 1985).

The arsenic forms associated with solid phases in aquatic sediments and their role in the biogeochemical cycling of As are largely unknown, at least partly because the relevant mineralogy has been mostly inferred from indirect measurements such as porewater concentrations of dissolved species and selective chemical dissolution of the solid sediment (Cullen and Reimer 1989; Maher 1984; Maher and Butler 1988; Woolson 1977). The precise mineralogy and crystal chemistry of these forms may play a significant role in determining bioavailability for many organisms (Lee et al. 2000; Maher 1984). In addition, environmental As/HFO materials are expected to be very similar (same column in the periodic table) in structure and chemistry to environmental P/HFO materials that are believed to play a dominant role in P cycles (Parfitt 1978), yet are easier to characterize by many physical methods that are sensitive to atomic mass (diffraction in particular). Greater knowledge of the precise mineralogy of key reactive phases (e.g., P/HFO and heavy metal HFO associations), in turn, would allow advanced reaction transport models of sediment diagenesis to be further developed (Van Cappellen and Wang 1996). Finally, As/HFO coprecipitation probably represents the main route of As removal from hydrothermal fluids and deposition to the sea floor, which is by far the largest geochemical reservoir of As (Maher and Butler 1988). These are some of the many reasons to obtain detailed mineralogical characterizations of natural As/HFO materials, that have motivated the present study.

Because of the importance of As/HFO associations in both natural and anthropogenic settings, there have been several detailed mineralogical studies of synthetic As-HFO sorbent-sorbant materials and coprecipitates (Raven et al. 1998; Wilkie and Hering 1996; Waychunas et al. 1996, 1995, 1993; Manseau 1995; Rea et al. 1994; Fuller et al. 1993; Pierce and Moore 1982). These studies have demonstrated the large As sorption capacity of HFO, showing that sorption densities as high as 0.7 and 0.25 mole-As/mole-Fe can be obtained in coprecipitation and sorption experiments, respectively (Fuller et al. 1993). Such demonstrated capacity far exceeds current technological limits using other sorbants, such as “granular ferric hydroxide” (a CI-bearing poorly crystallized β-FeOOH compound resembling akaganèite) which attains 0.06 mole-As/mole-Fe (Drieuhaus et al. 1998). The large demonstrated As sorption densities suggest that the As-HFO materials are intermediate between classic sorbate-sorbant systems and solid solutions or compounds in which the As and Fe occupy cation sites in a crystal structure that is more stable than a two-phase mixture. Indeed, the As is found to stabilize the HFO with respect to transformations to other oxides and to complex to the HFO surface in such a way as to modify the local coordination environments of a significant fraction of the Fe (Waychunas et al. 1996). The large sorption densities can be partly understood by the intrinsic nanocrystalline nature of 2-line ferrihydrite (typically 1–3 nm diameters) but the large values occurring in coprecipitation probably involve significantly smaller HFO particles (or less densely packed primary HFO particles) whose growth (or agglomeration) is poisoned by As complexation.

In this paper, we report a mineralogical characterization of the most As-rich (by two orders of magnitude) natural HFO precipitates ever described, that attain sorption densities of ~0.1–0.2 mole-As/mole-Fe. These As-sorbed HFO materials were formed by rapid abiotic precipitation when an As and Fe2+-rich hydrothermal fluid mixed with seawater, thereby preventing otherwise certain destruction of the surrounding coral reef ecosystem (Pichler et al. 1999a; Pichler and Veizer 1999; Pichler 1998).

Note that we follow current practice (Cornell and Schwertmann 1996) by reserving the term “2-line ferrihydrite” for the HFO material that has the corresponding characteristic powder X-ray diffraction (pXRD) pattern consisting of two well defined Bragg peaks. We use the term HFO for all 2-line ferrihydrite-like materials, including: (1) 2-line ferrihydrite itself, both in cases where it was identified as such and in cases where it is inferred to be 2-line ferrihydrite although the pXRD pattern was not measured; (2) materials that are less crystalline (or have smaller particle sizes) than 2-line ferrihydrite and that consequently have pXRD patterns that either do not have any pXRD peaks or have significantly broadened ones; and (3) sorbed or coprecipitated 2-line ferrihydrite-like materials having pXRD patterns that are qualitatively the same but may be significantly different in shape from the characteristic pXRD pattern of 2-line ferrihydrite. We take HFO materials to exclude the more crystalline ferrihydrite varieties (3-line to 6-line ferrihydrite).

SAMPLE PROVENANCE AND GEOLOGICAL SETTING

The collection area lies along the southwest margin of Ambitle Island (4°05’S, 153°40’W), one of the Feni islands in the southernmost island group of the Tabar-Feni chain, off the coast of Papua New Guinea. The island is part of a Quaternary stratovolcano with a central eroded caldera built on poorly exposed Oligocene marine limestone (Wallace et al. 1983). Volcanic strata (interbedded lava flows, lahar deposits, tuffs, and scoriae) dip radially from the island, presumably extending...
beneath the shelf. Several geothermal areas are located primarily along the western coast and in the western part of the caldera near breaches in the caldera wall.

Submarine hydrothermal venting occurs at Tutum Bay in shallow (5–10 m) water along the inner shelf that contains a patchy distribution of coral-algal reefs surrounded by medium to coarse-grained mixed carbonate-volcaniclastic sand and gravel (Pichler and Dix 1996). Two types of venting are observed: (1) focused discharge of a clear hydrothermal fluid occurs at discrete ports, 10–15 cm in diameter, where fluid temperatures at vent orifices are between 89 and 98 °C and (2) dispersed or diffuse discharge consisting of streams of gas bubbles emerging directly through the sandy to pebbly unconsolidated sediment and through fractures in volcanic rocks. A more detailed description of Tutum Bay is provided elsewhere (Pichler and Dix 1996, Pichler et al. 1999b) and color images can be found at http://chuma.cas.usf.edu/~pichler/.

As-bearing HFO materials are present throughout Tutum Bay where they form as thin layers on sediment grains in areas of high sea floor temperatures. Massive layers and extensive filling of sediment pore space also occurs but is restricted to the vicinity of vent sites where the HFO materials form bright orange coatings on volcanic boulders, distinct bands on coral skeletons and aragonite, and/or precipitate as massive layers in open spaces. The HFO precipitation at these vent sites has been studied by thermodynamic modeling and is found to occur as expected from such models, given the measured local physical and chemical conditions (Pichler 1998; Pichler and Veizer 1999; Pichler et al. 1999b). The HFO materials vary in color from a bright orange to a very dark brown that is almost black and have hardnesses, on Mohs’ hardness scale, from less than 1 to ~2.5. Banding into ~0.1–10 mm thick layers that are mostly parallel to the substrate surface and that have distinct color differences (predominantly yellowish vs. predominantly reddish) is common on all substrates. All samples were found to be stable in air after collection and were simply washed with deionized water and left to dry at room temperature before the various measurements were applied.

The samples studied in the present work are described in Table 1, in which their dry powder colors are described using the Munsell soil color chart system. Included are four relatively pure As-HFO samples (V2AR, V2BY, 97R, and 97Y) and one sample (FV1) that contained significant amounts of other mineral phases and that illustrates the evolution of the As-HFO materials toward other more stable Fe-bearing and As-bearing phases. Samples V2AR and 97R show intermediate degrees of evolution, with some crystalline phases being present. Such evolution has been described elsewhere and several more samples were studied by chemical analyses (Pichler 1998). Samples V2AR and V2BY are from the two main and adjacent layers (reddish and yellowish, respectively) that make up one deposit, as are samples 97R and 97Y (Table 1). These couples (V2AR/V2BY and 97R/97Y) were chosen in an attempt to quantify the physical and chemical differences that give rise to the color banding, as a first step in studying its mechanism and role in the overall aging of the samples. Note that the yellowish and reddish bands occur in reversed order in the two couples (Table 1). This may arise from the fact that the V2AR/V2BY couplet was in contact with seawater relatively far away from the vent, whereas the 97R/97Y couplet was directly exposed to the mixing hydrothermal fluid for some time (Table 1).

**EXPERIMENTAL AND ANALYSIS METHODS**

For transmission electron microscopy (TEM), dry samples were finely crushed in a glass mortar and dispersed in LR White embedding resin. Samples were twice infiltrated with the resin for 60 minutes before being re-suspended in fresh resin and polymerized at 60 °C for one hour. Thin sections (80–90 nm) were cut with a diamond knife mounted onto an ultramicrotome and collected on carbon coated Ni-grids having a supporting Formvar film. The thin sections were observed with a Philips EM400T electron microscope equipped with an X-ray spectrometer (LinK Analytical eXL/LZ-5) for elemental analysis. Energy dispersive spectroscopy (EDS) was performed at 100 kV with a beam current of 10 nA for ~100 seconds (live time). The size of the electron beam was 100 nm. Selected area electron diffraction (SAED) analyses were performed at 100 kV and the camera constant value was 640 nm.

The pXRD patterns were collected using an automated Philips X’Pert PW3710 system θ–2θ powder diffractometer, using CuKα radiation with a variable divergent slit and a solid state detector. Samples were mounted as acetone smears on a single-crystal Si wafer low-background holder that spun during data collection. The low-background holder gave a structur...
tureless background that cannot mask or interfere with the broad lines arising from the samples, in the entire 20 range of interest. Tests showed that control smears with only acetone also gave no detectable signals that could interfere with the signals from the samples. For pXRD and most measurements (vis-IR, Mössbauer, Munsell color determination, magnetometry, GC) the samples were not ground extensively but were powdered by simple tapping of dry pieces between sheets of wax paper, using a scalpel handle, in order to subject them to the least amount of mechanical treatment.

Room temperature \( (RT = 22 \, ^\circ C) \) \(^{57}\)Fe Mössbauer spectra were collected in transmission mode using random orientation powder absorbers having 27–41 mg of dry sample per cm² in holders with 1/2 inch diameter windows. Calibration was obtained using an enriched \(^{57}\)Fe foil at \( RT \) and all center shifts (CSs) are reported with respect to the CS of metallic iron at \( RT \). The transducer was operated in constant acceleration mode and folding was performed to achieve a flat background. All Mössbauer spectra were fit with the Voigt-based fitting method of Rancourt and Ping (1991) for quadrupole splitting distributions (QSDs), using the Recoil software of Lagarec and Rancourt (http://www.physics.uottawa.ca/~recoil/). Please consult the original article of Rancourt and Ping (1991) for all Mössbauer parameter definitions and a description of the relevant notation. All measurements were performed with a velocity range of ±11 mm/s and several measurements were performed on an expanded velocity range of ±4 mm/s. The larger range allowed possible detection of magnetically split patterns from bulk Fe oxides whereas the smaller range gave an optimal resolution for the As-HFO pattern. Fitting results for the two velocity ranges for a given sample were identical within statistical errors of the fitting parameters. Liquid nitrogen temperature \( (LNT = 77 \, ^\circ K) \) spectra were collected and calibrated in a similar way, in a cryostat in which the sample was also maintained at \( T = 77 \, ^\circ K \). As a result the CSs of the LNT spectra are shifted by an amount equal to the second order Doppler shift \( (SOD) \) of the source, that is, by the difference in SOD in going from \( RT \) to \( LNT \) [must add 0.12(1) mm/s to quoted \( LNT \) CS values].

The SQUID magnetometry measurements were performed on a low-field dc SQUID magnetometer that has been described elsewhere (Lamarche 1989). Typically a few milligrams of powder sample were immobilized in the cryogenic holder. The sample was first zero-field quenched (Earth’s field), then the field was turned on (typically 40–70 Gauss) and the magnetic moment on the sample was measured on warming (field warming, FW). After the maximum temperature was reached or at some other chosen temperature, one could choose to measure as the sample was recooled (field cooling, FC). The following field strengths were used for the various samples: 72.2 Gauss, samples 97R and 97Y, 52.2 Gauss, samples V2AR and V2BY; and 39.8, 40.7, and 40.0 Gauss for HFO samples synthesized with \( pH = 9, 7, \) and 6, respectively \((\text{Fe}) = 10^{-3} \, M) \).

Diffuse reflectivity spectra (vis-IR = 350–2100 nm) were obtained at \( RT \) with a Cary-14 spectrophotometer configured with a 23 cm diameter integrating sphere. The instrument has a double-beam double monochromator with a resolving power of 0.1–0.3 nm and a wavelength accuracy of ±0.4 nm. Sample and reference ports are located in the bottom of the sphere so that samples are measured in a horizontal orientation flush with port openings. Beam size (monochromatic light) at sample and reference ports was ±1.0 × 2.0 cm. Halon is used as the reflectivity standard and as the coating for the integrating sphere. Reflectivity measurements were converted to absolute reflectivity using literature data for the absolute reflectivity of Halon (Weidner and Hsia 1981). The compartment of the Cary-14 that contains the integrating sphere and samples was flushed with dry \( N_2 \) gas (derived from liquid \( N_2 \)) for at least 1 hour prior to obtaining reflectivity spectra. Reported spectra are the average of three scans, and the intensity of the bands at 1400 and 1900 nm, that result from \( H_2O \) and \( OH \), did not change during measurement. Because of limited amounts of sample (~25 to 160 mg), sample thickness was 0.012–0.080 g/cm² in sample cups that were Al metal painted with optically black paint. Assuming a powder density of 2.0 g/cm³, sample thickness is 0.06–0.40 mm. The work of Buckingham and Sommer (1983) suggests that a thickness of 0.025 mm is sufficient for conditions of effectively infinite thickness (i.e., incident light did not see the sample cup) for fine-grained powders of iron oxides.

XRF elemental analyses were obtained using a Philips PW2400/00 Sequential X-ray Spectrometer. Because of limited sample amounts (given the self-imposed constraint of using exactly the same sample fractions for all measurements), a special calibration was developed in which fused disks were prepared using only 20 mg of sample for 4.0750 g of flux material (3.2000 g LiB_4O_7; 0.8750 g LiBO_2). The calibration was for the eight most abundant elements excluding As. Extensive tests that involved repeated measurements on a given fused disk and comparisons with relevant standards that were not used in establishing the calibration allowed us to conclude that the accuracy of the measurement for the eight calibrated elements (Fe, Si, Na, Ca, Al, Mg, P, and K) was ~1 wt% and that the precision was ~0.1 wt%. The measured loss on ignition (LOI) of a given sample had an accuracy of ~3 wt% that was mainly determined by the uncertainty in weight loss of the flux, that is ~0.1 wt%. Arsenic was treated in a separate way, because of the difficulty in obtaining As-rich standards as fused disks. Background-corrected fluorescence peak counts were calibrated against a fused disk made using a synthetic Fe-As compound. The resulting accuracy in As is estimated at ~5%, including both uncertainties in the synthetic standard and systematic errors from not fully including matrix effects, whereas the comparative (sample to sample) precision is much better (~0.1%). The latter number is from the statistical error on the total number of background corrected counts for the given fluorescence peak. The magnitudes of As amounts found are in fair agreement with values for similar samples, obtained from neutron activation measurements (Pichler 1998).

Light element analyses (C, H, N, S) were obtained by gas chromatography (GC) of combustion products using thermal conductivity (thermistor) detection with a CE Instruments model EA-1110 elemental analyzer. Quantification was achieved by using time-integrated thermistor signals and corresponding run-specific calibrations for C, H, N, and S. Several standards were run before, during, and after the
measurements, that were used to construct a linear calibration curve for each element. Detection limit and accuracy are ~0.01 wt% of the element.

**Synthetic 2-line ferrihydrite reference material**

In order to correctly characterize our natural As-HFO samples and to elucidate the role of As in its crystal chemical structure and microstructure, it is important to have a relevant comparison point. The obvious choice is pure synthetic 2-line ferrihydrite, which is also formed by rapid precipitation and is believed to be the best candidate to model the sorbent phase of As-HFO, in a sorbate-sorbent picture.

In another study (unpublished) we have synthesized and extensively characterized several synthetic 2-line ferrihydrite samples, formed by rapid oxidation and precipitation of Fe⁺⁺ at pH values of 6, 7, 8, and 9, and at initial Fe concentrations of both 10⁻¹ or 10⁻⁴ M. In all cases, except at [Fe] = 10⁻³ M and pH = 8 where a mixture of lepidocrocite and goethite resulted, pure 2-line ferrihydrite was produced. The precipitates were washed with ultrapure water and freeze dried prior to measurements. All these synthetic ferrihydrite samples had measured characteristics (pXRD, Mössbauer spectra, TEM, compositions) similar to those for both natural and synthetic 2-line ferrihydrite (Jambor and Dutrizac 1998; Cornell and Schwertmann 1996; Childs 1992). They had variations of their characteristics on synthesis conditions (pH and [Fe]) that were much smaller than the differences in characteristics between any such synthetic or natural 2-line ferrihydrite sample and our As-HFO materials.

Therefore, although 2-line ferrihydrite is itself a complex material whose structure, microstructure, and crystal chemistry are far from well known (notwithstanding a recent high resolution TEM study that is of special note in this regard: Janney et al. 2000), it can be reproducibly synthesized and unambiguously identified. It also has sufficiently narrow spreads in measured characteristics, from one 2-line ferrihydrite sample to another, to make it an appropriate reference material for comparisons with its various sorbed, coprecipitated, or cation substituted varieties. Regarding Mössbauer spectroscopy, this is corroborated by a complete literature review of the RT parameters (CS and average QS) of natural and synthetic As-poor and As-free HFO materials (Thibault 2000). The latter study shows that the distributions of sample-specific average QS values for natural and synthetic samples are almost indistinguishable: mean = 0.75(2) mm/s, standard deviation = 0.062 mm/s, for natural samples (n = 10), compared to mean = 0.707(8) mm/s, standard deviation = 0.047 mm/s, for synthetic samples (n = 36), obtained from Lorentzian doublet fits. Both values of the mean are much smaller than typical values found for As-HFO samples. The slightly larger values of the mean and standard deviation for the natural samples may be due to the natural samples having various compositional differences, although the difference is small given the large variation in Si contents.

Indeed, a main complicating factor is the mostly unknown effects of Si in natural ferrihydrite materials, since our As-HFO samples do contain significant amounts of SiO₂. We expect the contribution from the effects of Si in our As-HFO samples to be small because unambiguous and systematic effects of Si on pXRD patterns and RT Mössbauer spectra of natural and synthetic ferrihydrite materials, with widely different Si/Fe ratios, have not been reported (Jambor and Dutrizac 1998; Childs 1992), except recently (e.g., Zhao et al. 1996; Glasauer et al. 2000). The latter authors convincingly argue that Si does not enter the HFO structure but instead complexes to the surfaces of HFO primary particles. Several results in the present study suggest that the same is true for As and that As and Si (or C) compete for reactive sites on HFO cores. In cases where authors report Si entering the HFO structure (e.g., Boyd and Scott 1999), the effects seen can also be attributed to other sample characteristics and they are small compared to the differences between As-HFO and As-free HFO reported here. In particular, the pXRD patterns of Si-rich HFO samples are not significantly different from those of Si-poor samples (Glasauer et al. 2000).

**Results and Discussion**

**Powder X-ray diffraction**

The pXRD patterns are shown in Figure 1 where they are compared to the pXRD pattern of a synthetic 2-line ferrihydrite ([Fe] = 10⁻³ M, pH = 9). Sample FV1 (Table 1) is an example of a presumably older deposit in which the As-HFO has largely transformed to a possibly As-bearing 6-line ferrihydrite, with its six main diffraction peaks labeled F, that coexists with several crystalline phases that give rise to prominent Bragg peaks that are labeled as follows: S for a smectite-illite mixture, B for bassanite (CaSO₄·0.5H₂O), and C for claudetite (As₂O₃). It is likely that all samples are predominantly As-HFO when freshly precipitated and that the solid phase composition of FV1 has arisen by aging in contact with seawater (e.g., Pichler 1998). In addition to bassanite, anhydrite (CaSO₄) was detected in one diffractogram of sample 97R where it may have resulted from post collection dehydration of a small amount of bassanite. The formation of bassanite (in FV1) is consistent with the known behavior of the aqueous CaSO₄·H₂O system, in the presence of NaCl and at the observed elevated sea floor temperatures (Deer et al. 1966). This would explain the relative hardness of FV1 (Table 1) in terms of filing of sediment pore space by authogenic bassanite.

The reddish samples (V2AR and 97R) are predominantly As-HFO and contain small amounts of smectite-illite mixtures and claudetite. The yellowish samples (V2BY and 97Y) are predominantly As-HFO with no detectable amounts of other solid phases, except for halite (NaCl) that appears in some diffractograms and that is probably produced from residual seawater during drying. That the amounts of crystalline phases in the reddish samples (V2AR and 97R) are small is clear from: (1) reference intensity ratio calculations applied to the diffraction peaks; (2) our TEM investigation of the same samples (described below); and (3) our bulk chemical analyses (e.g., total Ca, described below). Nonetheless, the reddish samples seem to have been more susceptible than the yellowish samples to the type of aging that presumably has occurred in FV1. The initial aging sequence appears to be from As-HFO (2-line As-
bearing HFO) to less As-rich 6-line ferrihydrite with excess Si and some Fe forming a smectite-illite mixture and excess As forming claudetite, accompanied by consolidation via authigenic bassanite. Smectite-illite mixtures are common in hydrothermal environments (e.g., Ylagan et al. 2000) and these T-O-T type layer silicate mixtures may initially form from the Fe³⁺ layered octahedra of HFO having tetrahedral Si⁴⁺ complexed to their faces. The 2-line to 6-line evolution of the predominant As-HFO component is unusual in that such evolution has not been reported in other natural environments (Cornell and Schwertmann 1996) and may be made possible in part by the elevated sea floor temperatures in our case. In terrestrial surface environments and in related laboratory experiments HFO evolves directly to hematite or goethite (Cornell and Schwertmann 1996).

Overall, on first examination, the pXRD patterns of the As-HFO samples (Fig. 1) show two very broad Bragg humps at d spacings of approximately 1.5 and 2.5 Å which are generally taken to be characteristic of 2-line ferrihydrite (thus the name). This might bring one to conclude that the samples are simply 2-line ferrihydrite. In fact, more information can be derived from the pXRD patterns. Such pXRD patterns are directly related to the underlying radial distribution functions (i.e., angle-averaged Patterson functions that include all disorder and finite particle size effects), as is well known in the study of amorphous materials and liquids. This means that the differences in pXRD pattern shapes translate into differences in populations of cation-cation, cation-anion, and to a lesser extent anion-anion distances. This was used by Waychunas et al. (1996) to study synthetic As-HFO materials, in comparison with extended X-ray absorption fine structure (EXAFS) results. In light of the latter study, the systematic differences in the 1.5 and 2.5 Å peaks of our As-HFO samples and 2-line ferrihydrite (Fig. 2), consisting in more intensity on the high d spacing sides of the As-HFO peaks, can be interpreted as being due to a significant population of Fe-As structural pairs, across shared bridging coordination O anions on the surface (Waychunas et al. 1996). The presence of a significant population of nearest neighbor As-O pairs can also be inferred and the general smearing out (broadening), relative to the Bragg peaks of 2-line ferrihydrite, is understood in terms of smaller X-ray coherent domains (i.e., smaller primary particles) in the As-HFO materials.

One concludes from the pXRD measurements (Figs. 1 and 2) that our As-HFO materials are significantly different from As-free 2-line ferrihydrite and that they have patterns like those of synthetic As-HFO samples (Waychunas et al. 1996). The As is not in a separate crystalline phase (the very small amounts of claudetite in the reddish samples notwithstanding) and there is evidence for Fe-As pairs that preclude most of the As from being in separate amorphous or nanocrystalline phases. The pXRD patterns are correlated to sample color in that the reddish samples have small amounts of secondary phases but the latter secondary phases (mainly smectite-illite) would be ex-
expected to make the samples less red by reducing the relative amounts of iron oxyhydroxides. The yellowish color must be attributed to As complexation in As-HFO, in samples (V2BY and 97Y) that do not contain detectable secondary phases, as corroborated by several other measurements described below.

Transmission electron microscopy

TEM micrographs for samples 97Y and 97R are shown in Figure 3. Such micrographs, combined with extensive TEM-based SAED and EDS measurements, allow us to conclude: (1) that there is no evidence for separate phases on this size scale or for phases other than the predominant As-HFO seen by pXRD; (2) that there is no evidence for bacteria or associated biologically induced mineralization, which is common in other hydrothermal vent environments (Fortin et al. 1998; Rona 1988; Juniper and Fouquet 1988); and (3) that the As-HFO particles are agglomerated into large spheroidal objects having typical dimensions 100–1000 nm. Chemical analyses for samples 97Y and 97R based on an average of seven EDS probe spots, in areas where the beam size was smaller than the sample structure, are given in Table 2 as elemental ratios, where they are compared to the bulk XRF values for all As-HFO samples. The TEM-based EDS analyses are semiquantitative in that we expect particle morphology and differences in compactness to significantly affect the results.

Samples 97Y and 97R display somewhat different submicrometer-scale morphologies (Fig. 3), with spheroidal and elongated agglomerations typically having dimensions 150–250 nm (diameters of spheroids), 200 × 300 nm, 400 × 500 nm, and 600 × 1000 nm in 97Y and 100–150 nm (diameters of spheroids) and 150 × 350 nm in 97R. The agglomerations are composed of nanometer-scale As-HFO particles that are detected as dense spots in Figure 3 (arrows) and that are typical of ferrihydrite particles seen at this magnification.

Chemical analyses: XRF and GC

The XRF results (Table 3) clearly demonstrate a systematic difference between the reddish samples (97R and V2AR) and the associated yellowish samples (97Y and V2BY, respectively) in which the reddish samples have: less Fe₂O₃, less As₂O₅, and more SiO₂. This translates (Table 2) into the reddish samples having significantly smaller As/Fe and larger Si/Fe ratios. These differences are not seen by the TEM-based EDS method because of its semiquantitative nature with such samples. Such systematic compositional differences between reddish and yellowish As-HFO samples from this geothermal area have not previously been reported but they are real. They are also important, as they give a chemical basis for the observed differences in color and other physical properties, such as the susceptibility to smectite-illite formation described above and the Mössbauer parameters and magnetism described below. They suggest that there is competition between As and Si for surface complexation to HFO. A similar competition, between P and Si, has been observed in sorption/desorption experiments with ferrihydrite (Parfitt 1989) and another similar competition, between C (as sorbed onto HFO from natural organic matter) and As, was observed in a study of As removal from drinking water (Hering et al. 1997).

The only other example that we are aware of where As-HFO material forms banded structures with alternating reddish and yellowish layers is the case described by Leblanc et al. (1996) of ferruginous accretions produced in an acid mine water environment. The latter material is similar to our samples in many respects. Their As contents and layer thicknesses are comparable to ours and, as is the case with our samples, the reddish layers are harder, have more advanced crystallization to more stable phases, and contain significantly less As. Their materials, however, do not contain significant amounts of Si.

Apart from the red/yellow differences, the chemical analysis results (Tables 2 and 3) are in general agreement with previous evaluations on related samples (Pichler 1998; Pichler and Veizer 1999; Pichler et al. 1999a) and show the samples to be very As-rich (0.1–0.2 mole As/mole Fe) but to otherwise be similar to natural As-free ferrihydrite samples (Jambor and Dutrizac 1998), in their Si and minor element contents and in their LOI values. Natural and synthetic (Saleh and Jones 1984; Vempati et al. 1990) ferrihydrite samples can typically contain such large Si/Fe ratios as those found in our As-HFO samples. There have been no systematic studies of the effect of Si on the pXRD patterns of 2-line ferrihydrite samples and there also have been no conclusive reports that Si has a systematic effect on the pXRD patterns of ferrihydrite samples of any type (Jambor and Dutrizac 1998). Indeed, many Si-rich 2-line ferrihydrite samples are found to have pXRD that are not significantly different from those of Si-poor 2-line ferrihydrite samples (e.g., Glasauer et al. 2000; their Fig. 1 compared to our Figs. 1 and 2). Also, the large differences in SiO₂ contents...
**Figure 3.** TEM thin-section micrographs and sample EDS spectra for As-HFO samples: (a) 97Y and (b) 97R. The arrows point to examples of primary particles, that are seen as dark ~1–5 nm spots at this magnification. The Ni in the EDS spectra is from the Ni grid sample supports that are used. The CuKα line has an interference with the NiKβ line, such that Cu cannot reliably be quantified but in any case occurs in relatively small amount.
TABLE 2. Molar ratios from TEM-based EDS analyses for 97R and 97Y, compared to values from XRF analyses

<table>
<thead>
<tr>
<th>Sample</th>
<th>As/Fe</th>
<th>Si/Fe</th>
<th>Ca/Fe</th>
</tr>
</thead>
<tbody>
<tr>
<td>97R</td>
<td>0.13</td>
<td>0.26</td>
<td>0.035</td>
</tr>
<tr>
<td>97Y</td>
<td>0.12</td>
<td>0.34</td>
<td>0.034</td>
</tr>
</tbody>
</table>

† pH = 7; [Fe] = 10^-3 M; freeze dried.

* Using only 97R-1 from Table 3, because 97R-2 seems anomalous in both As and LOI.

The amounts of H reported in Table 4 are consistent with ~1 nm diameter primary particles having all their surface cations (Fe^3+, As^5+, Si^4+, C^4+) being coordinated on the free surface side by OH– and OH_2⁻. This is true for both As-HFO and synthetic HFO.

Table 3. XRF results* (in wt% of assumed oxide)

<table>
<thead>
<tr>
<th>Sample</th>
<th>Fe_2O_3</th>
<th>As_2O_3</th>
<th>SiO_2</th>
<th>Na_2O</th>
<th>CaO</th>
<th>Al_2O_3</th>
<th>MgO</th>
<th>P_2O_5</th>
<th>K_2O</th>
<th>LOI</th>
<th>Total</th>
</tr>
</thead>
<tbody>
<tr>
<td>97R-1</td>
<td>52.1</td>
<td>13.6</td>
<td>21.9</td>
<td>3.6</td>
<td>1.9</td>
<td>1.9</td>
<td>1.6</td>
<td>0.3</td>
<td>0.3</td>
<td>15.9</td>
<td>113.1</td>
</tr>
<tr>
<td>97R-2</td>
<td>50.0</td>
<td>4.9</td>
<td>19.1</td>
<td>2.1</td>
<td>1.7</td>
<td>0.9</td>
<td>1.4</td>
<td>0.2</td>
<td>0.3</td>
<td>27.9</td>
<td>107.9</td>
</tr>
<tr>
<td>97Y-1</td>
<td>57.8</td>
<td>22.1</td>
<td>12.4</td>
<td>3.7</td>
<td>2.0</td>
<td>0.9</td>
<td>1.2</td>
<td>0.3</td>
<td>0.2</td>
<td>18.4</td>
<td>119.0</td>
</tr>
<tr>
<td>97Y-2</td>
<td>56.7</td>
<td>18.8</td>
<td>11.4</td>
<td>1.1</td>
<td>1.8</td>
<td>0.0</td>
<td>1.1</td>
<td>0.3</td>
<td>0.2</td>
<td>22.4</td>
<td>113.8</td>
</tr>
<tr>
<td>97Y-3</td>
<td>53.4</td>
<td>20.0</td>
<td>10.7</td>
<td>1.4</td>
<td>1.7</td>
<td>0.0</td>
<td>1.1</td>
<td>0.3</td>
<td>0.2</td>
<td>10.9</td>
<td>99.7</td>
</tr>
<tr>
<td>97Y-4</td>
<td>53.2</td>
<td>16.4</td>
<td>10.7</td>
<td>1.7</td>
<td>1.8</td>
<td>0.0</td>
<td>1.1</td>
<td>0.3</td>
<td>0.2</td>
<td>13.9</td>
<td>99.3</td>
</tr>
<tr>
<td>97Y-5</td>
<td>56.9</td>
<td>16.2</td>
<td>11.0</td>
<td>1.6</td>
<td>1.8</td>
<td>0.0</td>
<td>1.0</td>
<td>0.3</td>
<td>0.2</td>
<td>12.9</td>
<td>101.9</td>
</tr>
<tr>
<td>V2AR</td>
<td>52.6</td>
<td>10.7</td>
<td>17.4</td>
<td>2.7</td>
<td>2.2</td>
<td>1.0</td>
<td>1.0</td>
<td>0.4</td>
<td>0.2</td>
<td>17.9</td>
<td>106.1</td>
</tr>
<tr>
<td>V2BY</td>
<td>59.5</td>
<td>19.1</td>
<td>14.9</td>
<td>2.4</td>
<td>2.3</td>
<td>1.1</td>
<td>0.9</td>
<td>0.2</td>
<td>0.2</td>
<td>16.4</td>
<td>117.0</td>
</tr>
<tr>
<td>FV1</td>
<td>62.4</td>
<td>13.0</td>
<td>20.3</td>
<td>4.3</td>
<td>2.0</td>
<td>0.5</td>
<td>1.4</td>
<td>0.2</td>
<td>0.2</td>
<td>14.9</td>
<td>119.2</td>
</tr>
<tr>
<td>syn-HFO†</td>
<td>85.1</td>
<td>0.0</td>
<td>0.0</td>
<td>1.8</td>
<td>0.0</td>
<td>0.0</td>
<td>0.0</td>
<td>0.0</td>
<td>0.0</td>
<td>13.9</td>
<td>100.8</td>
</tr>
</tbody>
</table>

* All other elements analyzed (Ce, Co, Cr, Cu, Mn, Ni, Pb, Ti, V, Zn) were relatively minor components (less than ~0.1 wt%). Large totals for As-bearing samples are explained in the text.
† pH = 7; [Fe] = 10^-3 M; freeze dried.

REPORTED TEXT:

...allows one to conclude that the observed different susceptibilities to form (Si-bearing) smectite-illite phases and claudetite (Fig. 1). For these reasons, and because Si has a smaller atomic number than Fe and As, we did not consider the presence of Si in our above interpretation of the pXRD results. The same holds for the presence of C, that we have quantified by GC as described below.

The GC light element analysis results are given in Table 4. It is interesting to note a significant amount of C that is correlated to sample color in the same way as Si (Table 2). The reddish samples (97R and V2AR) have approximately twice as much C as their yellowish counterparts (97Y and V2BY, respectively). This C is probably not organic (TEM results) and may be complexed to the HFO surfaces in much the same way as the Si. As noted above, the measured H contents allow a quantitative interpretation of the measured LOI, that also holds for synthetic HFO. Interestingly, the weight fractions of H for As-HFO and for synthetic HFO are very similar (1.8–2.1 wt%) yet when expressed as an H/Fe molar ratio there is a significant difference (factor of ~2). This suggests that OH and H_2O yet when expressed as an H/Fe molar ratio there is a significant difference (factor of ~2). This suggests that OH and H_2O are not primarily coordinating Fe in the HFO structure but rather coordinating the surface cations indiscriminately, in a roughly constant weight fraction because of roughly constant specific surface areas of reactive surfaces. In other words, it appears that, irrespective of the amounts of Si, C, and As that complex to the surfaces of HFO primary particles, there is always approximately the same number of required OH or H_2O anions per weight of sample to complete the inner sphere coordinations of the exposed or outermost cations. Indeed, a simple calculation that uses typical cation-oxygen bond lengths and the most common coordination geometries allows one to conclude that...
form of a bending or stretching vibration involving As, that the As-HFO samples contain significant amounts of As. Certain features may be due to the effects of As but these are difficult to identify because there have not been systematic studies of the effects of As on the vis-IR diffuse reflectivity spectra of Fe oxides or oxyhydroxides. The low wavelength region up to ~1 µm is well understood in terms of ligand field transitions of Fe³⁺ in octahedral coordination (Scheinost et al. 1998; Clark et al. 1990). The broad absorption band at ~1.93 µm is due to sorbed and coordinating H₂O whereas the band at ~1.43 µm is due to both H₂O and OH groups (Clark et al. 1990). The spectra therefore offer proof of significant amounts of H₂O, in both As-HFO and synthetic HFO. The presence of significant amounts of OH cannot be ruled out or conclusively established, based on the vis-IR spectra alone.

There are four main features that differentiate the spectra for As-HFO from that for synthetic HFO. First, the overall magnitude of reflectivity is significantly smaller for synthetic HFO, in the entire wavelength region. This is also true when synthetic HFO is compared to natural ferrihydrite (Cornell and Schwertmann 1996) and may be due to an absence of spectroscopically inert (in this region) SiO₂ in the synthetic sample. Second, the main visible-region absorption edge is at ~0.65 µm for synthetic HFO, compared to ~0.58–0.63 µm for As-HFO. The latter absorption edge is known to correlate well with Munsell hue, as is the case here (2.5YR for synthetic HFO, 5YR for 97R, Table 1). We provide an explanation for the latter shift in the next section (Mössbauer section). Third, the H₂O + OH band minimum is at ~1.40 µm for synthetic HFO whereas it is at ~1.43 µm for As-HFO and for natural ferrihydrite (Cornell and Schwertmann 1996). This band is also relatively narrow in synthetic HFO. These facts may be interpreted as evidence for synthetic HFO containing more OH and less H₂O, as is also suggested by its relatively shallow 1.93 µm band. Alternatively or in addition to this, the OH groups in synthetic HFO will coordinate Fe whereas they are likely to coordinate surface complexed As, Si, and C in As-HFO. This may shift the OH stretch frequency by an amount sufficient to explain the observed difference. Finally, the small wavelength edge of the 1.93 µm absorption band is at a smaller wavelength in synthetic HFO than in As-HFO. Given the large width of this band, this may be an artifact of saturation effects.

**57Fe Mössbauer spectroscopy**

A typical fitted RT Mössbauer spectrum (sample 97R) is shown in Figure 5, where it is compared to the RT Mössbauer spectrum of a typical synthetic 2-line ferrihydrite sample (pH = 9, [Fe] = 10⁻³ M). All the spectra of our As-HFO samples (97R, 97Y, V2AR, V2BY, and FV1), on both velocity ranges used (±11 and ±4 mm/s), are very similar to each other and do not show any detectable amounts of other Fe-bearing phases, such as the common bulk magnetic Fe-oxides (e.g., hematite, goethite) that would show easily detectable sextets (e.g., Kodama et al. 1977). There is a large difference between the RT Mössbauer spectra of our As-HFO samples and the RT Mössbauer spectra of all known ferrihydrite and ferrihydrite-like materials, including our synthetic 2-line ferrihydrite samples (Thibault 2001), all previously reported synthetic and natural ferrihydrite materials (Cardile 1988; Childs and Johnston 1980; Coey and Readman 1973b; Henmi et al. 1980; Johnston and Lewis 1983; Kauffman and Hazel 1975; Loseva and Murashko 1973; Madsen et al. 1986; Mathalene et al. 1970; Mizota and Maeda 1986; Murad 1988a; Murad 1989; Murad et al. 1988; Murad and Schwertmann 1980; 1988; Rea et al. 1994; Saraswat et al. 1977; Schwertmann and Murad 1988; Van Der Giessen 1967; Vandenberghe et al. 1990; Winnik et al. 1998), ferritin (Brady et al. 1968), and schwertmanite (Murad 1988b). The LNT spectra of ferritin are similar to those expected at RT (St-Pierre et al. 1986). All the latter materials (irrespective of Si contents) have RT Mössbauer spectra that are relatively similar to each other but very different from the RT Mössbauer spectra of our As-HFO samples. This difference is clearly seen in the spectra themselves (Fig. 5) and translates into significantly different fitting parameters, on further analysis.

**Figure 4.** Vis-IR diffuse reflectivity for four As-HFO samples and one synthetic HFO sample (pH = 7, [Fe] = 10⁻³), as indicated.

**Figure 5.** RT Mössbauer spectrum of As-HFO sample 97R, compared to that of a synthetic HFO sample ([Fe] = 10⁻³ M, pH = 9). The fit to the 97R spectrum, that is described in Tables 3 and 4, is shown by a solid line running through the data points and the corresponding residual is shown above the spectra, using the same vertical scale. The two spectra have been rescaled and shifted to facilitate comparison.
Unfortunately, most previous studies have fitted the *RT* spectra of ferrihydrite with Lorentzian doublets, most often using either one or two doublets per spectrum. While this approach provides a quantitative assessment of the spectra, it has many problems such as non-unique solutions arising from parameter tradeoffs and lack of physical meaning of the individual doublets. Rancourt (1994) has shown that this approach is incorrect and has described the problems associated with it. The correct approach in this application is to acknowledge that the physical cause of the absorption lineshape is inhomogeneous broadening arising from a distribution of local environments that gives rise to a QSD (Murad 1989; Vandenbergh et al. 1990). The QSD represents the most crystal chemical information that can be extracted from a spectrum (Rancourt et al. 1994a; Rancourt 1998).

The fit parameters for both *RT* and *LNT* spectra are given in Table 5. Those samples run only at ±11 mm/s (FV1, V2AR, V2BY) were fit with only two QSD components and a CS-QS coupling parameter (δ) that was confined to be zero. This gave greater fitting stability and avoided certain artifacts such as extended high-QS QSD tails. Extensive tests showed that, barring such artifacts, the same QSDs were obtained with either two or three (or more) QSD components, within statistical error. Also, freezing δ at zero or letting it adjust had no noticeable affect on the extracted QSD. The spectra collected at *LNT* (samples 97R and 97Y) were fit with two Gaussian QSD components and δ was allowed to have a non-zero value. This was sufficient to give statistically acceptable fits in these cases. Finally, we note that all fits of *RT* spectra, of both synthetic HFO and the As-HFO samples, were significantly improved if we allowed ~5% of intrinsic absorption area asymmetry (A−/A+, ~1.05) and that such allowed asymmetry did not change any of the extracted QSDs. Such asymmetry and its physical causes has been described previously in a nanophase oxide (Rancourt et al. 1985) but has not previously been reported in ferrihydrite. It will not be considered further in this article.

Associated calculated parameters are given in Table 6 and are defined as follows: <QS> is the average magnitude (i.e., average absolute value) of the quadrupole splitting (QS), σQSD is the standard deviation of the QSD, and <CS> is the average CS. These fits and the resulting parameters show that all detectable Fe in our As-HFO samples (and in all ferrihydrite-like

### Table 5. Mössbauer fit parameters

<table>
<thead>
<tr>
<th>Sample/temperature</th>
<th>BG (MC/ch)</th>
<th>Area (kC.mm/s)</th>
<th>δ0 (mm/s)</th>
<th>δ1 (%)</th>
<th>pL (mm/s)</th>
<th>Δl (mm/s)</th>
<th>σQSD (mm/s)</th>
<th>χ2/ν</th>
</tr>
</thead>
<tbody>
<tr>
<td>97R RT</td>
<td>1.0735</td>
<td>0.2716</td>
<td>0.362</td>
<td>-0.0016</td>
<td>5</td>
<td>1.6</td>
<td>1.1</td>
<td></td>
</tr>
<tr>
<td>97Y RT</td>
<td>3.16955</td>
<td>0.7214</td>
<td>0.3703</td>
<td>-0.0096</td>
<td>6</td>
<td>1.9</td>
<td>0.86</td>
<td>1.58</td>
</tr>
<tr>
<td>V2AR RT</td>
<td>3.90153</td>
<td>1.2578</td>
<td>0.3617</td>
<td>0.0†</td>
<td>43.7</td>
<td>1.17</td>
<td>0.483</td>
<td>5.82</td>
</tr>
<tr>
<td>V2BY RT</td>
<td>2.72229</td>
<td>1.0235</td>
<td>0.3604</td>
<td>0.0†</td>
<td>22.5</td>
<td>1.45</td>
<td>0.610</td>
<td>11.9</td>
</tr>
<tr>
<td>FV1 RT</td>
<td>3.13538</td>
<td>0.9350</td>
<td>0.3603</td>
<td>0.0†</td>
<td>16.7</td>
<td>1.42</td>
<td>0.72</td>
<td>12.9</td>
</tr>
<tr>
<td>syn.-HFO‡ RT</td>
<td>0.84436</td>
<td>0.22779</td>
<td>0.3494</td>
<td>-0.0032</td>
<td>14</td>
<td>1.1</td>
<td>0.53</td>
<td>7.92</td>
</tr>
<tr>
<td>97R LNT</td>
<td>0.13746</td>
<td>41.910</td>
<td>0.3435</td>
<td>-0.0040</td>
<td>32</td>
<td>0.589</td>
<td>0.29</td>
<td>0.61</td>
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<tr>
<td>97Y LNT</td>
<td>0.107084</td>
<td>31.690</td>
<td>0.3478</td>
<td>-0.0069</td>
<td>49</td>
<td>0.735</td>
<td>0.35</td>
<td>0.72</td>
</tr>
</tbody>
</table>

* All parameters are as defined by Rancourt and Ping (1991). In all fits, the Lorentzian full width at half maximum, Γ, was frozen at the known Heisenberg value of 0.097 mm/s. This allows QSDs to be compared directly. Also, only symmetric elemental doublets were used, with A−/A+ = 1.0, where the latter ratio is the ratio of the spectral area of the low energy Lorentzian line in the elemental doublet to the spectral area of the high energy Lorentzian line in the elemental doublet. Either two or three QSD Gaussian components were used, depending on the number required to obtain a satisfactory fit. BG is the flat background level, in megacounts per channel. Area is the total spectral area. δ0 and δ1 are the linear coupling coefficients between CS and QS, as: CS = δ0 + δ1(QS). pL is the percent probability of the i-th Gaussian component of the QSD. Δl and σQSD are the center and sigma-width, respectively, of the i-th Gaussian component of the QSD.

† This parameter was frozen during the fit.

‡ pH = 9; [Fe] = 10⁻³ M; freeze dried.

### Table 6. Calculated Mössbauer parameters

<table>
<thead>
<tr>
<th>Sample</th>
<th>Temperature</th>
<th>&lt;CS&gt;* (mm/s)</th>
<th>&lt;QS&gt; (mm/s)</th>
<th>σQSD (mm/s)</th>
<th>&lt;QS&gt; (mm/s)</th>
<th>Munsell hue (YR)</th>
</tr>
</thead>
<tbody>
<tr>
<td>97R RT</td>
<td>0.3607</td>
<td>0.913</td>
<td>0.506</td>
<td>0.859</td>
<td>5</td>
<td></td>
</tr>
<tr>
<td>97Y RT</td>
<td>0.3608</td>
<td>0.986</td>
<td>0.499</td>
<td>0.957</td>
<td>10</td>
<td></td>
</tr>
<tr>
<td>V2AR RT</td>
<td>0.3617</td>
<td>0.918</td>
<td>0.494</td>
<td>0.890</td>
<td>7.5</td>
<td></td>
</tr>
<tr>
<td>V2BY RT</td>
<td>0.3604</td>
<td>0.899</td>
<td>0.503</td>
<td>0.962</td>
<td>10</td>
<td></td>
</tr>
<tr>
<td>FV1 RT</td>
<td>0.3603</td>
<td>0.826</td>
<td>0.504</td>
<td>0.852</td>
<td>5</td>
<td></td>
</tr>
<tr>
<td>syn.-HFO‡ RT</td>
<td>0.3470</td>
<td>0.744</td>
<td>0.373</td>
<td>0.717</td>
<td>2.5</td>
<td></td>
</tr>
</tbody>
</table>

* With respect to the CS of metallic iron at *RT*, except in LNT measurements where one must add 0.12(1) mm/s (as explained in text). <CS>, <QS>, and σQSD are from the QSD fits (Table 5) whereas <QS>, is the average QS obtained from Lorentzian doublet fits, using 2 or 3 Lorentzian doublets arranged symmetrically about similar values of CS.

† pH = 9; [Fe] = 10⁻³ M; freeze dried.
materials) is Fe$^{3+}$, having parameters consistent with octahedral O coordination, and that there is no detectable tetrahedral Fe$^{3+}$. Tetrahedrally coordinated Fe$^{3+}$ would have a significantly lower CS (Coe 1980; Rancourt et al. 1992) and would give rise to a $^{10}$Fe$^{3+}$ QSD centered at a smaller average QS value and having a smaller width than the $^{10}$Fe$^{3+}$ QSD (Rancourt et al. 1992). Also, the QSs and CSs of $^{16}$Fe$^{3+}$ and $^{10}$Fe$^{3+}$ are expected to have significantly different temperature dependencies (Rancourt et al. 1994b) that would give rise to significant changes in spectral shape on going to LNT, for example. This does not occur. The $RT$ and $LNT$ spectra for a given sample are very similar to each other and have very similar QSD shapes (Tables 5–6). Our conclusion is in contradiction to Cardile’s (1988) position that significant amounts of tetrahedral Fe$^{3+}$ may be consistent with the $RT$ Mössbauer spectra of ferrihydrite.

The QSDs extracted from the $RT$ Mössbauer spectra of all our As-HFO samples and of synthetic 2-line ferrihydrite (pH = 9, $[\text{Fe}] = 10^{-3} M$) are compared in Figure 6. As is the case with the spectra themselves, the QSDs of all the As-HFO samples are similar to one another. The most different one from the group (which also has the lowest $\langle QS \rangle$, Table 6) is sample FV1. This sample is also the oldest sample that has lost much of its As to a separate As-oxide phase (Fig. 1 and Pichler 1998; Pichler and Veizer 1999). The QSDs show the large systematic differences between the $RT$ Mössbauer spectra of our As-HFO samples and As-free ferrihydrite, directly in terms of the local environments of Fe, as expressed in the QSs. The As-HFO samples have significant populations of local environments with large QSs (~1.5 mm/s) whereas As-free ferrihydrite samples do not. The As-HFO QSDs are centered on larger $\langle QS \rangle$ values, are broader, and have their maxima (i.e., most probable QS values) at ~0.85 mm/s instead of at the typical value of ~0.60 mm/s for As-free ferrihydrite samples.

The QSDs shown in Figure 6 suggest that a large fraction of the Fe (~20%) has local environments that are significantly perturbed by the direct presence of As, via shared coordinating O anions. Given our measured As/Fe ratios (Table 2), this fraction is consistent with a very small primary particle size, such as the sections of double octahedral chains containing only 4–10 Fe$^{3+}$ cations proposed by Waychunas et al. (1996), keeping in mind that one As atom can significantly distort two Fe octahedra in a bidentate arsenate sorption geometry such as the one proposed by Waychunas et al. (1995; 1996). Note that a primary particle containing ~10 Fe atoms and the corresponding numbers of As, Si, and C atoms, consistent with our chemical analyses, would be ~1 nm in diameter, which is comparable to the diameters of primary particles recently seen in synthetic 2-line ferrihydrite by high resolution TEM (Janney et al. 2000). The fact that ~20% of the Fe has local environments that are significantly perturbed by the direct presence of As also means that ~80% of the Fe has local environments similar to those in As-free HFO and allows us to conclude that the As and Fe probably do not mix and are separated in structures consisting of Fe-rich primary particle cores having surface complexed As. This, in turn, suggests that a relatively stable and long lived nucleus of HFO must first be formed before As can complex to it and that Fe-As complexes do not initiate viable nucleus formation during precipitation. The latter scenario is consistent with the Si-As and C-As competitions for surface complexation sites on HFO cores discussed above.

Although the differences between the QSDs of the relatively pure As-HFO samples (97R, 97Y, V2AR, V2BY; Fig. 6) are small compared to the difference between the QSD of synthetic HFO and the QSDs of the As-HFO samples, they are significant. This is seen on close examination of the corresponding $\langle QS \rangle$ values (Table 6). The average QS is known to be a robust property of the QSD, that can be estimated from the most common Lorentzian doublet fits if care is taken to avoid CS-QS inter-doublet tradeoffs. Our $\langle QS \rangle$ values are compared (Table 6) to $\langle QS \rangle$ values, $\langle QS \rangle_{T}$, obtained by taking area weighted averages from Lorentzian doublet fits. The latter estimates are seen to systematically underestimate the average QS by ~0.02–0.06 mm/s (0.035 mm/s average underestimate). This arises because the QSD measure is a true average of the QS magnitude (since the QSD is a distribution of QS magnitudes) whereas in a Lorentzian doublet approach the effectively negative QS contributions (that have no physical meaning in a paramagnetic doublet context) are not properly taken into account. Nonetheless, $\langle QS \rangle_{T}$ is a useful quantity because it can be extracted from the most commonly reported Lorentzian doublet fits, thereby allowing important comparisons with published data. As noted above, for example, the $\langle QS \rangle_{T}$ values for synthetic As-free 2-line ferrihydrite samples occur in the 1σ range 0.66–0.75 mm/s (36 samples) and the $\langle QS \rangle_{T}$ values for natural As-poor HFO samples occur in the 1σ range 0.69–0.81 mm/s (10 samples) (Thibault 2001). The differences in $\langle QS \rangle$ (or in $\langle QS \rangle_{T}$ from one As-HFO sample to another (and with the synthetic HFO sample and literature values) are much larger than the accuracy in determining $\langle QS \rangle$. Let us therefore next examine the $\langle QS \rangle$ values in more detail.

A most striking feature (Table 6) is that $\langle QS \rangle$ is strongly correlated to the Munsell hue (Table 1 and repeated in Table 6 for convenience). Without exception, there is a near linear relation between the Munsell YR index values (0 is most reddish, 10 is most yellowish) of the As-HFO samples and their $\langle QS \rangle$ values. A strong correlation is maintained if one adds...
the synthetic HFO comparison point, which has the smallest value of Munsell hue YR index. The reddish samples (97R and V2AR) have significantly smaller <QS> values than their yellowish counterparts (97Y and V2BY, respectively), in proportion to their Munsell hue YR index values. This is most probably not an accident and there are good physical reasons to expect such a correlation. The QS is a measure of the degree of local distortion in the octahedral coordination environment of Fe\(^{3+}\) and of the lattice contribution to the local electric field gradient and such local distortions and field gradients also determine the molecular orbital energy levels of the Fe\(^{3+}\) cation, which in turn determine the main absorption band edge position in the visible, which is the main spectral feature that is associated with the Munsell hue (Scheinost et al. 1998). This correlation therefore establishes that the colors (hues) of our As-HFO samples are mainly determined by the Fe\(^{3+}\) cations in their local environments and not by any other chemical constituents (As, Si, C, ...), except indirectly if the latter constituents cause distortions of the Fe\(^{3+}\) environments. Since Si and C do not cause significant distortions of the Fe\(^{3+}\) environments (as measured by Mössbauer QSs) and since As content is strongly correlated to Munsell hue (Tables 1–3), As must be the main cause of the increased QSs and associated changes in the Fe\(^{3+}\) local environments. There can be no doubt that As is in close coordination proximity to Fe. The quite striking yellowish hues of our As-HFO samples are caused by As, which, by distortion, distorts the local environments of a large fraction of the Fe\(^{3+}\) cations thereby changing their electronic orbital structures in such a way as to cause a significant shift of the main visible absorption band edge. It is likely that this mechanism also causes the yellowish colors that are correlated to As contents in Si-poor As-rich As-HFO formed in acid mine drainage environments (Leblanc et al. 1996), in further support of the conclusion that Si does not play a significant role in determining the hue. We also note that the small amounts of smectite-illite that are detected by pXRD in the reddish samples (V2AR and 97R) are not expected to contribute significant features to the measured QSDs. Indeed, the known <QS>\(_2\), of the main Fe\(^{3+}\)-bearing smectite, nontronite, is 0.44(4) mm/s (average for 22 samples, unpublished), a value that precludes any interference with the main QSD feature at ~1.5 mm/s that we attribute to As complexation.

Finally, our <QS> values for As-HFO samples measured at RT (Table 6) can be compared to two relevant measurements. First, Rea et al. (1994) measured the <QS>\(_2\), (calculated from a Lorentzian doublet fit) of a synthetic ferrihydrite sample in which the most labile surface sites were isotopically enriched in \(^{57}\)Fe and obtained 0.87 mm/s (their sample MF6). This value is significantly larger than typical values for ferrihydrite materials and approaches or surpasses our values for As-HFO samples. For example, it surpasses the <QS>, value of 0.829 mm/s for sample FV1 which has presumably lost much of its As to a more stable As-oxide phase. From this comparison we conclude that both the most labile sites in As-free samples and As-perturbed sites in As-HFO have larger QSs (and local distortions) than interior sites but that the perturbing effect of As is much larger than that of the surface. Second, Rea et al. (1994) also measured the <QS>, of a synthetic sample that was first surface enriched in \(^{57}\)Fe and then sorbed with As and obtained 0.78 mm/s. The latter value is in the range of As-poor natural HFO samples (0.69–0.81 mm/s) and is close to the 1σ range of As-free synthetic HFO samples (0.66–0.75 mm/s). It is also lower than their value for As-free labile sites and this was interpreted to mean that As decreased the distortion of the surface sites. The latter result and interpretation is in conflict with our results that demonstrate a large <QS>\(_1\) enhancement arising from As. This may be due to a fundamental structural difference between samples sorbed with As after synthesis vs. samples that are the result of rapid coprecipitation or it may be due to transformations that occurred under the particular isotopic exchange and sorption conditions that were used (Rea et al. 1994).

**SQUID Magnetometry**

Magnetometry is a powerful materials characterization tool and is particularly useful with iron oxides and oxyhydroxides because they exhibit various magnetic behaviors, such as magnetic ordering, spin-flop transitions, electron localization, mixed valences, superparamagnetism, superferromagnetism, and incomplete antiferromagnetic sublattice cancellation, that are very sensitive to their crystal chemical and microstructural features (e.g., Dang et al. 1998; Dunlop and Özdemir 1997). Magnetism has often been used as a sensitive characterization of HFO-like materials, including synthetic and natural ferrihydrite samples, a commercial ferrihydride-like catalysis material (NANOCA T, manufactured by Mach I, Inc.), and various types of ferritin, using both variable field and temperature Mössbauer spectroscopy (Bauminger and Nowik 1989; Cianchi et al. 1994; 1992; Coey and Readman 1973a; Hawkins et al. 1994; Madsen et al. 1986; Murad et al. 1988; Pankhurst and Pollard 1992; Pollard et al. 1992; Quin et al. 1988; St-Pierre et al. 1996; 1989; Zhao et al. 1996) and classic magnetometry measurements (Allen et al. 1998; Coey and Readman 1973b; Gider et al. 1995; Ibrahim et al. 1994; 1995; Lear 1987; Van Der Giessen 1967).

The zero-field quench FW susceptibilities (M/H) for our As-HFO samples (excluding FV1 which has significant amounts of crystalline phase impurities) and for three synthetic HFO samples are shown in Figures 7 and 8, respectively, where they are compared to the theoretical curves for paramagnets containing the same numbers of ferric cations per gram as synthetic HFO and natural Si-bearing As-HFO. The theoretical susceptibility of a paramagnet is given by the well known Curie law, in which we assume orbital quenching (L = 0) and a total electronic spin S = 5/2 for ferric iron. The only required sample-dependent parameter is the number of paramagnetic cations (Fe\(^{3+}\)) per gram of sample. All other ionic and molecular species (OH\(^–\), O\(^2–\), As\(^{3+}\), As\(^{5+}\), Si\(^4+\), Na\(^+\), Ca\(^2+\), Al\(^3+\), P\(^5+\), K\(^+\), Mg\(^2+\), C\(^6+\), and H\(_2\)O) are strictly diamagnetic. Assuming any of the most common stoichiometric formulas (such as Fe\(_2\)O\(_3\)-4H\(_2\)O, 5Fe\(_2\)O\(_3\)-9H\(_2\)O, or Fe\(_2\)O\(_3\)-2FeOOH·2.6H\(_2\)O; see Jambor and Dutrizac 1998) for ferrihydrite leads to the same number of Fe per unit mass: N\(_g\) = 6.27 × 10\(^{21}\) Fe/g. Using As/Fe = 0.15 and Si/Fe = 0.3 as representative numbers (Tables 2 and 3 and associated discussions) and ignoring C (Table 4), we then assume a reasonable formula for As-HFO of 5Fe\(_2\)O\(_3\)-9H\(_2\)O·3SiO\(_2\)-0.75As\(_2\)O\(_3\) to obtain N\(_g\) = 4.585 × 10\(^{21}\) Fe/
A superparamagnetism to superferromagnetism transition occurs at 25 K in these samples. The FC values are always larger than or equal to the FW values. FC data merge at ~25 K for sample V2B and at ~45 K for sample V2AR. The FW and FC data are shown for samples V2AR and V2B. The FW and FC data are shown for the pure paramagnetism of a representative As-HFO (dashed line). Both paramagnetism of a representative synthetic HFO (solid line) and to the theoretical susceptibilities corresponding to the pure paramagnetism of a representative synthetic HFO (solid line) and to the pure paramagnetism of a representative As-HFO (dashed line). Only FW data is shown for pH = 6 and 7, whereas both FW and FC data is shown for pH = 9. The superparamagnetism to superferromagnetism transition occurs at 55–59 K in these samples.

The resulting theoretical susceptibilities are shown as solid and dashed curves, respectively, in Figures 7 and 8. These curves give a reference point corresponding to a total absence of inter-moment interactions, such as those that cause superparamagnetism, superferromagnetism, antiferromagnetism and ferrimagnetism, and they allow a clearer comparison of HFO and As-HFO.

Superparamagnetism is the phenomenon whereby the supermoments on individual primary particles, that arise from intra-particle inter-atomic magnetic exchange interactions at temperatures sufficiently below the intrinsic magnetic ordering temperature (Néel or Curie point) of the material, have random fluctuations in orientation due to thermal excitation, much as individual atomic moments have thermal fluctuations in the usual paramagnetic state (Dormann et al. 1997). Superferromagnetism is the phenomenon in which the supermoments of superparamagnetic particles order magnetically, under the influence of inter-particle magnetic interactions. The latter interactions can be purely dipolar in nature or can have an exchange component from inter-particle bridges. The unusual magnetism of ferrihydrite and HFO samples has been interpreted in terms of these phenomena (Bauminger and Nowik 1989; Cianchi et al. 1994; 1992; Hawkins et al. 1994; Madsen et al. 1986; Murad et al. 1988; Pankhurst and Pollard 1992; Pollard et al. 1992; Quin et al. 1988; St-Pierre et al. 1996; 1989; Allen et al. 1998; Coey and Readman 1973a; 1973b; Gider et al. 1995; Ibrahim et al. 1994; 1995; Van Der Giessen 1967; Zhao et al. 1996), where the primary (~1–5 nm) particles have a transition from paramagnetism to superparamagnetism, as temperature is decreased, at or near the intrinsic magnetic ordering temperature of the material (ferrihydrite), \(T_{SPSF}\) followed by a transition from superparamagnetism to superferromagnetism at some lower temperature, \(T_{SPSF}\). The blocking temperature, \(T_B\), is defined as the temperature at which, on decreasing the temperature, the superparamagnetic fluctuations are observed to effectively stop, in a given measurement. Clearly, \(T_B\) is not a uniquely defined property of the sample in that it depends on the measurement time of the probe being used. The relation between \(T_B\) and \(T_{SPSF}\) and relevant concepts have been discussed by Rancourt and Daniels (1984). Because of the relation to various technological applications, the area of small particle magnetism is presently one of intense fundamental research, with significant advances in understanding being made in both superparamagnetism (e.g., Coffey et al. 1998) and superferromagnetism (e.g., Garcia-Otero et al. 2000; Nowak et al. 2000).

The magnetic behaviors of our As-HFO and synthetic HFO samples can be understood in terms of the above concepts and the usual ferrimagnetism that is common in iron oxide particles (e.g., Dang et al. 1998). The main features and their interpretation are described below. We find that the magnetic behaviors of HFO and As-HFO are dramatically different, far more than the mass dilution of paramagnetic cations caused by the adsorbed arsenate and incorporated or adsorbed Si or C would suggest. The main differences are related to significant differences in structure and microstructure, as discussed below.

Superparamagnetic/superferromagnetic transitions are seen in all the samples, where \(T_{SPSF}\) is estimated as the temperature at which a sharp cusp (or maximum) occurs in the FW curve. The latter temperature of the maximum in the FW curve is also usually close to the temperature at which temperature hysteresis sets in. That is, the temperature at which the FW and FC curves depart significantly on cooling. In all our synthetic HFO samples (pH = 6, 7, 9; [Fe] = 10^{-3} M), we find \(T_{SPSF} = 55–59\) K (Fig. 8) whereas all the As-HFO samples (except FV1 which...
has significant amounts of crystalline phase impurities; Table 1, Fig. 1) have clear signs of a transition at $T_{SP/SF} = 25 \pm 2$ K (Fig. 7). In the yellowish (more As-rich, 97Y and V2BY) samples a true FW maximum occurs at 25 K whereas in the reddish (97R and V2AR) samples a weaker peak is seen at 25 K to be superposed on a paramagnetic or superparamagnetic Curie-like contribution (Fig. 7). This suggests that the reddish samples may have inhomogeneous magnetic responses with some of the material in a given reddish sample having a superferromagnetic ordering temperature of 25 K and the other fraction of the material having lower superferromagnetic ordering temperatures, below 4.2 K since another FW peak is not seen. The Curie-like contributions in the FW and FC curves of the reddish samples may be superparamagnetic signals from fractions within the samples that have low superferromagnetic ordering temperatures. Part of the Curie-like contributions may also arise from the ordinary paramagnetism of exchange isolated or loose Fe$^{3+}$ moments.

$T_{SP/SF}$ is a consequence of several factors such as: supermoment magnitude, $\mu$, mean inter-particle distance, $d$, strength of inter-particle magnetic bridging, via superexchange shared oxygen bonds, etc. For example, for pure dipolar inter-particle interactions and a given primary particle spatial configuration, we have $T_{SP/SF} \propto \mu^2/d^6$. If primary particle packing is optimum and the supermoment is proportional to the primary particle volume, then $T_{SP/SF} \propto r^3$, where $r$ is the primary particle radius. This suggests that As-HFO material having $T_{SP/SF} \sim 25$ K has primary particles that are simply smaller than the synthetic HFO primary particles, by $\sim 25\%$. As-HFO material having $T_{SP/SF}$ approximately $< 4$ K would have primary particles that are much smaller than the primary particles in synthetic HFO. Alternatively, if the inter-particle separation is the same for As-HFO and synthetic HFO, one can interpret the lower value of $T_{SP/SF} = 25$ K for As-HFO as meaning that the material having this lower value of $T_{SP/SF}$ has supermoments that are $\sim 30\%$ smaller than the supermoms of synthetic HFO. Similarly, $T_{SP/SF}$ approximately $< 4$ K values would imply a much smaller supermoment magnitude. On the other hand, inter-particle exchange interaction bridging may play a role and sorbed OH/OH$_2$ coordinated complexes of As (and Si and C, to different extents) may poison such bridging in the As-HFO agglomerations of primary particles, thereby causing the observed reduced values of $T_{SP/SF}$.

The susceptibilities themselves have typical magnitudes that are much greater for our synthetic HFO samples than for our As-HFO samples. There is approximately a factor of 10 difference in the susceptibility magnitudes (Figs. 7–8), with the HFO susceptibilities being significantly larger than expected paramagnetic HFO values (Fig. 8) and the As-HFO susceptibilities being significantly smaller than expected paramagnetic As-HFO values (Fig. 7). This unusual situation can be understood in terms of certain features of small magnetic particles, as follows. The ratio of susceptibilities for ordinary paramagnetic (Curie law) behavior and for superparamagnetic (Langevin-type) behavior can be expressed as the ratio of the usual Curie constant, $C_J$, for cations of spin $J$, and the effective Curie constant, $C_L$, for Langevin-type superparamagnetism. Assuming one supermoment per primary particle, it can be shown that $C_J/C_L = (m^2/\pi^2) [J/(J+1)]$ (1) where $n$ is the number of moment-bearing cations in a primary particle and $m$ is the number of atomic moments per supermoment. One atomic (or cationic) moment has magnitude $\mu = g\mu_B$ where $g$ is the Landé $g$-factor and $\mu_B$ is the Bohr magneton. In all cases considered here $J = S = 5/2$ for Fe$^{3+}$ but the factor $m^2/\pi^2$ in Equation 1 can be: (a) comparable to 1, if $m \sim n^{1/2}$ as expected for random vacancies in a two-sublattice antiferromagnet (see Dormann et al. 1997); (b) larger than 1, if $m > n^{1/2}$ as occurs with ferromagnetic particles where $m^2/n \propto r^1$; or (c) smaller than 1, if $m < n^{1/2}$ as occurs when, for example, a sufficiently small antiferromagnetic particle has mostly paired opposing nearest neighbor moments. Clearly, in small particle systems all situations are possible. Our results suggest that the synthetic HFO samples and the As-HFO samples differ significantly in their mechanisms for creating a resulting net supermoment per primary particle. Synthetic HFO must have $m^2/n > 1$ whereas As-HFO must have $m^2/n < 1$. There is also a systematic difference between the susceptibility magnitudes of the reddish (97R and V2AR) and yellowish (97Y and V2BY) samples (Fig. 7) that implies that the yellowish (As-rich) samples have values of $m^2/n$ that are smaller than those of the reddish samples. Differences in $m^2/n$ can arise from purely geometric effects related to primary particle size and shape but can also involve differences in intra-particle surface moment coupling due to the presence or absence of sorbed As. The above Mössbauer results show that a significant fraction of the Fe in As-HFO has significantly distorted local environments, relative to those in HFO. Such distortions in Fe-O bond lengths and bond angles are known to significantly affect inter-atomic magnetic superexchange bond strengths. This is consistent with the more As-rich (yellowish) samples having the smallest values of $m^2/n$, the order being synthetic-HFO > reddish-As-HFO > yellowish-As-HFO.

On examination of the magnetometry results for the As-HFO samples (Fig. 7), one notes that the susceptibilities for samples 97Y and V2BY are almost identical, whereas they are step-wise different from the susceptibilities for samples 97R and V2AR, which are very similar to each other. This grouping correlates with the colors of the samples (Table 1), their chemical compositions (Tables 2–4), and their Mössbauer parameters (Tables 5–6). The measurements taken together, and the magnetometry results in particular, suggest a bimodal population of As-HFO deposits consisting of materials that are either reddish or yellowish in character, where As-HFO materials of intermediate character are avoided. In other words, the observed banding (Picher 1998; Pichler and Veizer 1999) seems to be well defined.

**FURTHER DISCUSSION AND CONCLUSION**

We find the natural As-HFO to be very similar to those synthetic As-HFO materials formed by coprecipitation, as opposed to sorption of As after HFO synthesis, and to be significantly different from all As-free related materials, including all known synthetic and natural HFO materials and ferrihydrite samples. As is the case for 2-line ferrihydrite and most HFO materials, a low degree of crystallinity (including both structural disorder
and small particle size effects) does not allow a structural determination but the data does permit many features of the microstructures and local structures to be described. All the data either support or are consistent with As-HFO being a nanophase material in which there is intimate molecular-scale association of tetrahedrally coordinated arsenate with Fe in a typical oxihydroxide framework of octahedrally coordinated Fe \(^{3+}\) and in which there is significant inter-particle agglomeration (or bridging or consolidation). The As seems to prevent particle growth and cause smaller HFO primary particles (or cores) to be formed, even compared to the effects of surface complexed Si and C. It also causes more structural disorder, in the form of significantly distorted octahedral environments of Fe. The latter distortions that are caused by As, since Si and C have much weaker effects, perturb the \(^{57}\)Fe \(^{3+}\) electronic orbitals and their energy levels which, in turn, results in both larger QSs and shifted visible absorption band edges, thereby causing the observed strong correlation between Munsell hue and Mössbauer parameters. In addition, the QS of Fe \(^{3+}\) is expected to have a large lattice (i.e., ionic point charge distribution) contribution that will also be affected by the presence of nearby As \(^{3+}\). The As also appears to be relatively strongly complexed to the HFO core, compared to Si or C, such that it more effectively retards evolution toward more stable crystalline phases that, initially, are smectite-illite mixtures and claudelite.

Overall, the observed effects that can be unambiguously attributed to As, as opposed to other sorbed or surface complexed species such as Si and C, include: (1) modifications of the pXRD pattern, relative to the 2-line patterns of As-free or As-poor HFO materials, indicating the intimate association of Fe and As; (2) relatively large QS values in the RT Mössbauer spectra that correspond to significantly distorted octahedral Fe \(^{3+}\) environments, again attesting to an atomic scale association of Fe and As; (3) a shift toward smaller wavelengths of the main absorption band edge in the visible and a corresponding increase of the Munsell hue YR index that occur in the natural As-HFO materials toward the more stable phases, including the initial stages that involve the distinct yellowish (samples 97Y and V2BY) and reddish (samples 97R and V2AR) layers that we have described. The type of As–Si/C avoidance interaction described above may drive the banding that occurs in the As-HFO deposits, via post-deposition surface diffusion across compact agglomerations of primary particles. Alternatively, the banding may be a record of variations in hydrothermal fluid or seawater characteristics. More work is also needed to more fully characterize the As-HFO materials themselves.

High resolution TEM, along the lines of the recent work on synthetic HFO by Janney et al. (2000), would be particularly welcome in the latter area, as would the use of ab initio quantum mechanical calculations of sorption and complexation, along the lines of the recent work by Randall et al. (1999).

Acknowledgments

Financial support from the Natural Sciences and Engineering Research Council of Canada is gratefully acknowledged. Thomas Pichler acknowledges support in the form of Geological Society of America student research grants. We thank Nathalie Roux, Milena Kusnir, and Gabriella Giustiniano for performing several of the syntheses and characterizations related to the synthetic ferrihydrite reference material and co-investigator Denis Mavrocordatos for allowing us to use unpublished data from the later study. The TEM/EDS measurements were performed at the University of Guelph, in T.J. Beveridge’s laboratory. This work was supported in part by NASA’s Cosmochemistry Program under RTOP 344-31-20-24 (R.V.M.). We thank Ron Hartree for his help in establishing the required XRF calibrations and procedures. We thank Mei-Zhen Dang, Wendy Abdi, and Gilles St-Jean for their work in establishing the required GC calibrations and procedures. We have benefitted from the often very detailed and extensive comments of four anonymous reviewers whom we thank.

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